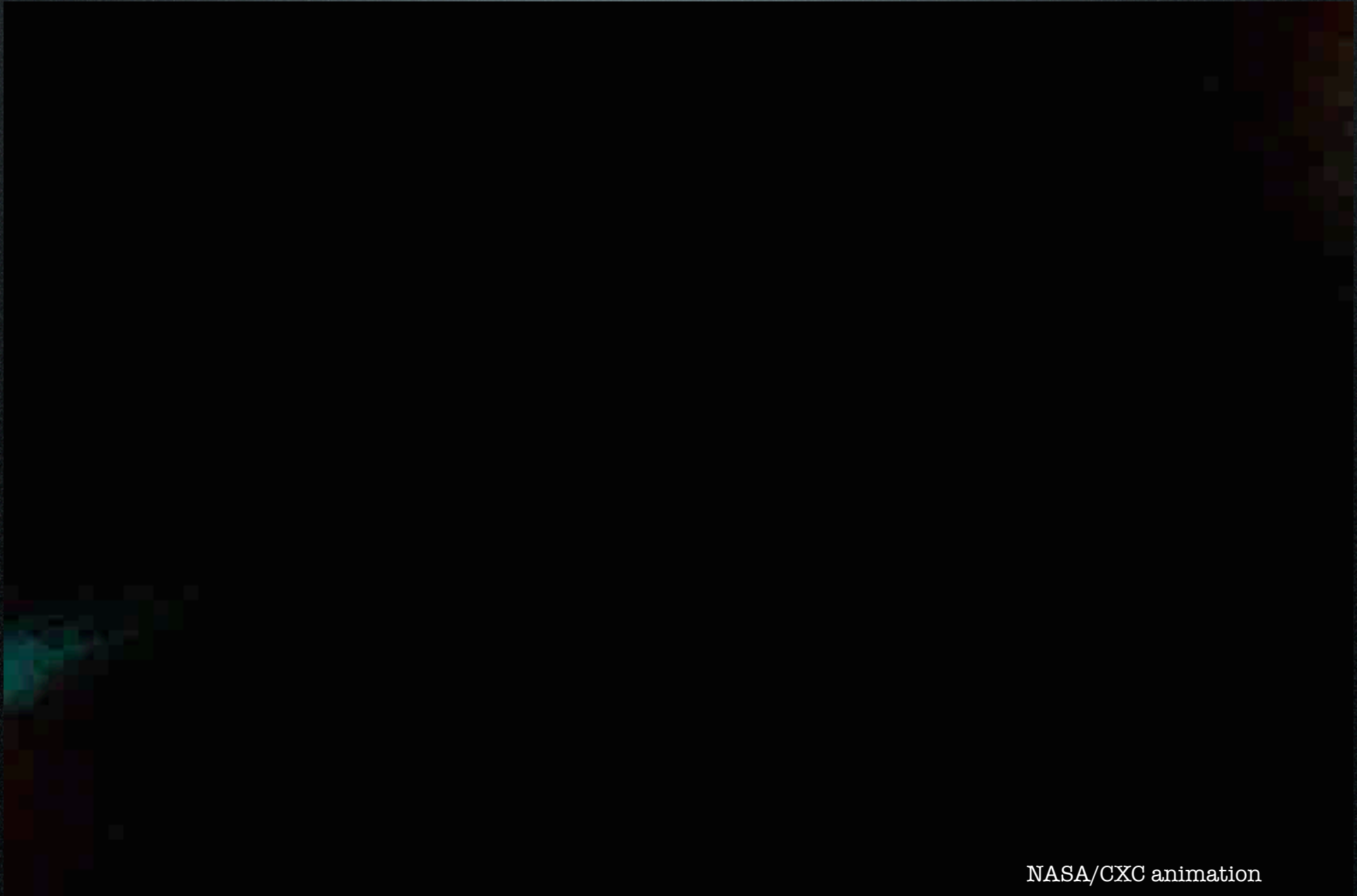


Supermassive black hole hierarchical evolution

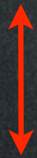


Outline

- 1. SMBHs in the local universe: where from?**
- 2. SMBHs Mass Growth: Accretion vs Merging**
 - AGN at low redshift
- 3. Dynamical Evolution of MBHs:**
 - Formation and coalescence of MBH binaries
 - Gravitational rocket
- 4. Consequences on MBHs spin**
- 5. Gravitational waves**

M_{BH} - σ relation: co-evolution of SMBHs and galaxies

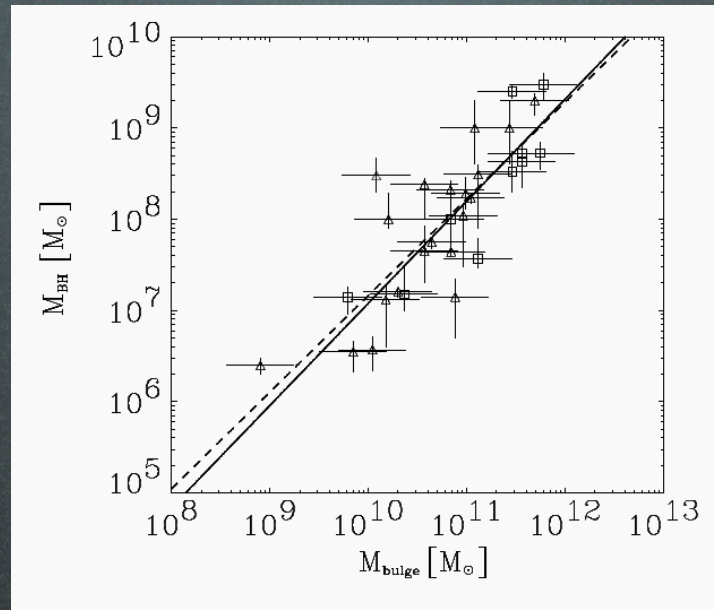
SMBHs



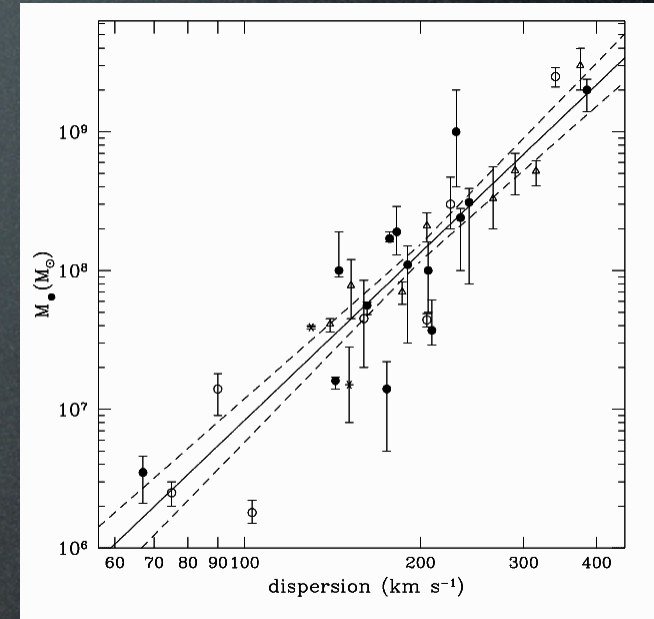
stellar bulges



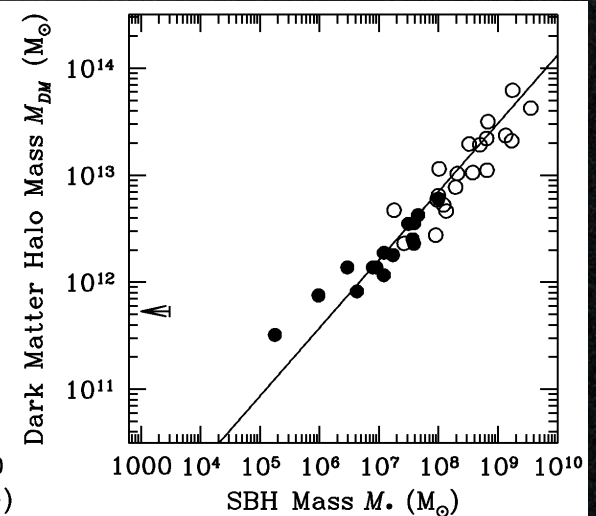
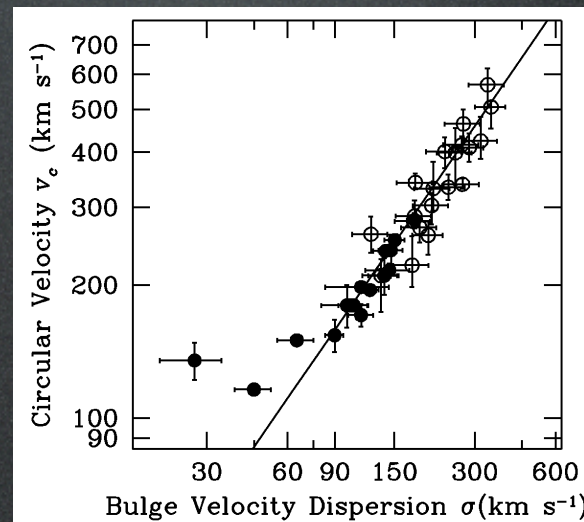
host DM halos (???)



from Marconi & Hunt 2004



from Tremaine 2002

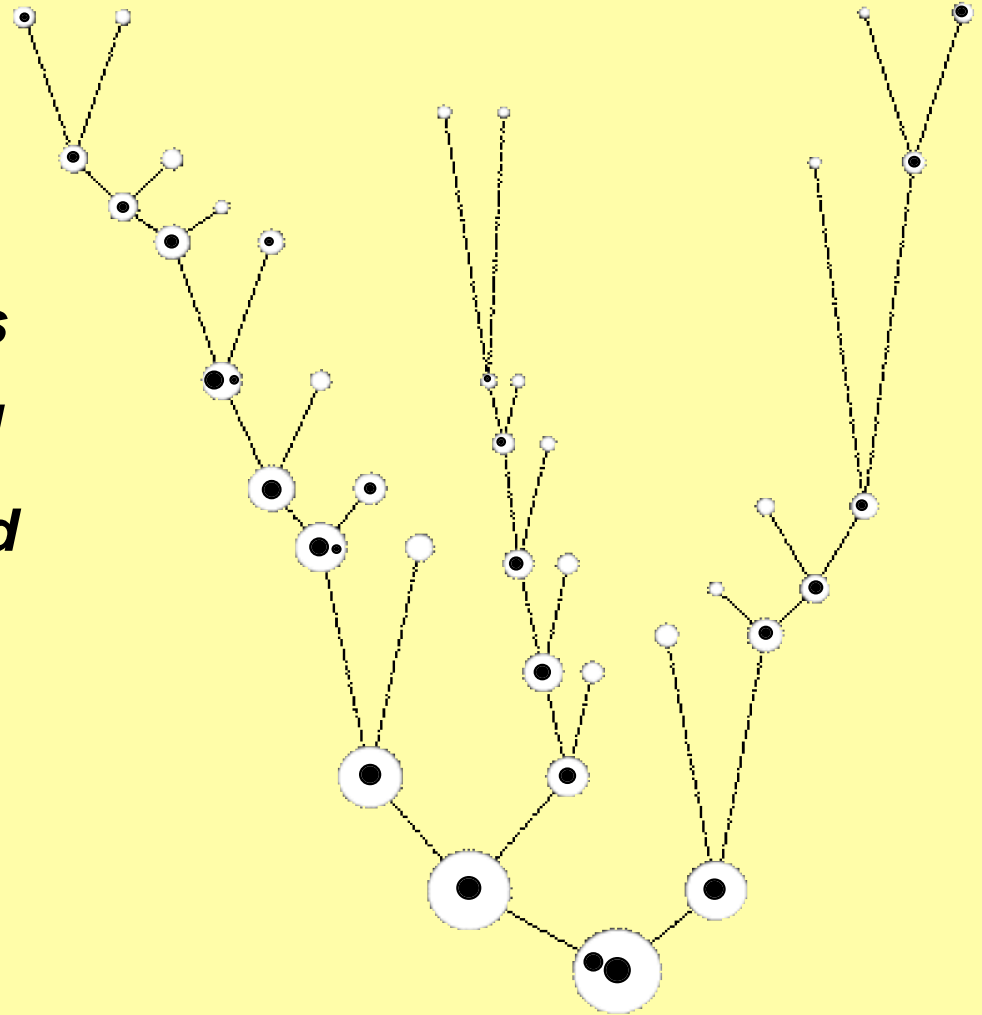


from Ferrarese 2002

THE MODEL

SMBHS are grown from *seed* pregalactic BHs. These seeds are incorporated in larger and larger halos, *accreting gas* and *dynamically interacting* after mergers.

Note: in most of what follows I'll consider seeds which are the endproduct of the *first stars*. Results at low-*z* are unchanged



Volonteri, Haardt & Madau 2003

Chapter 2

The seeds at $z > 20$ are tiny, $\sim 100-10^4 M_{\text{sun}}$

How do they grow to become supermassive?

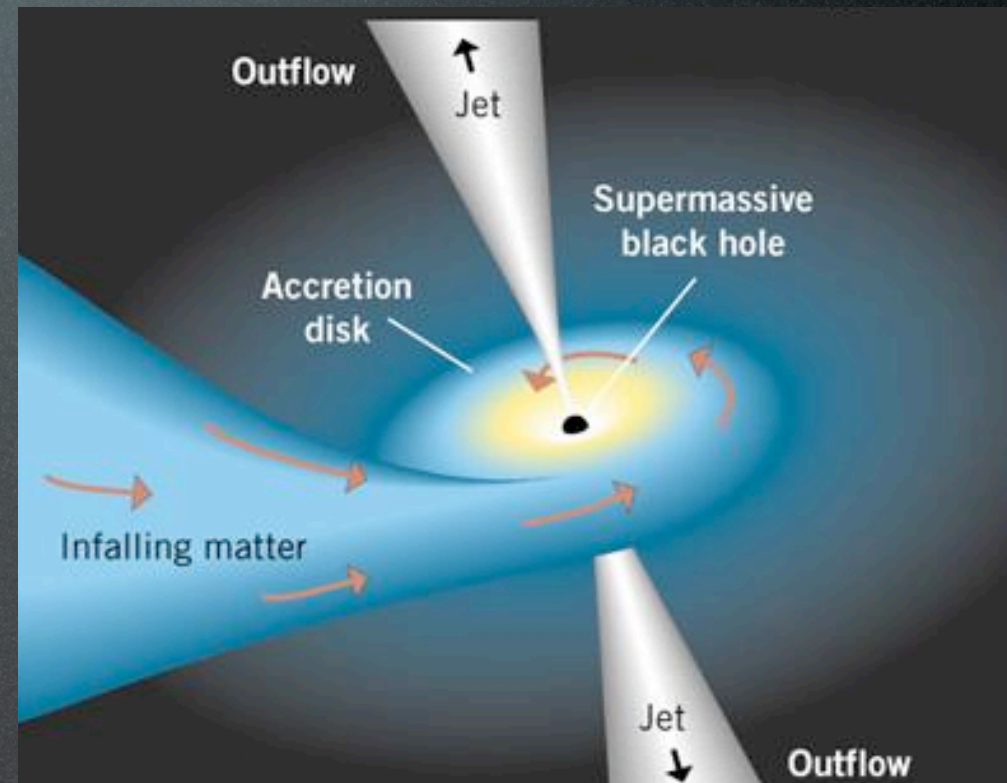
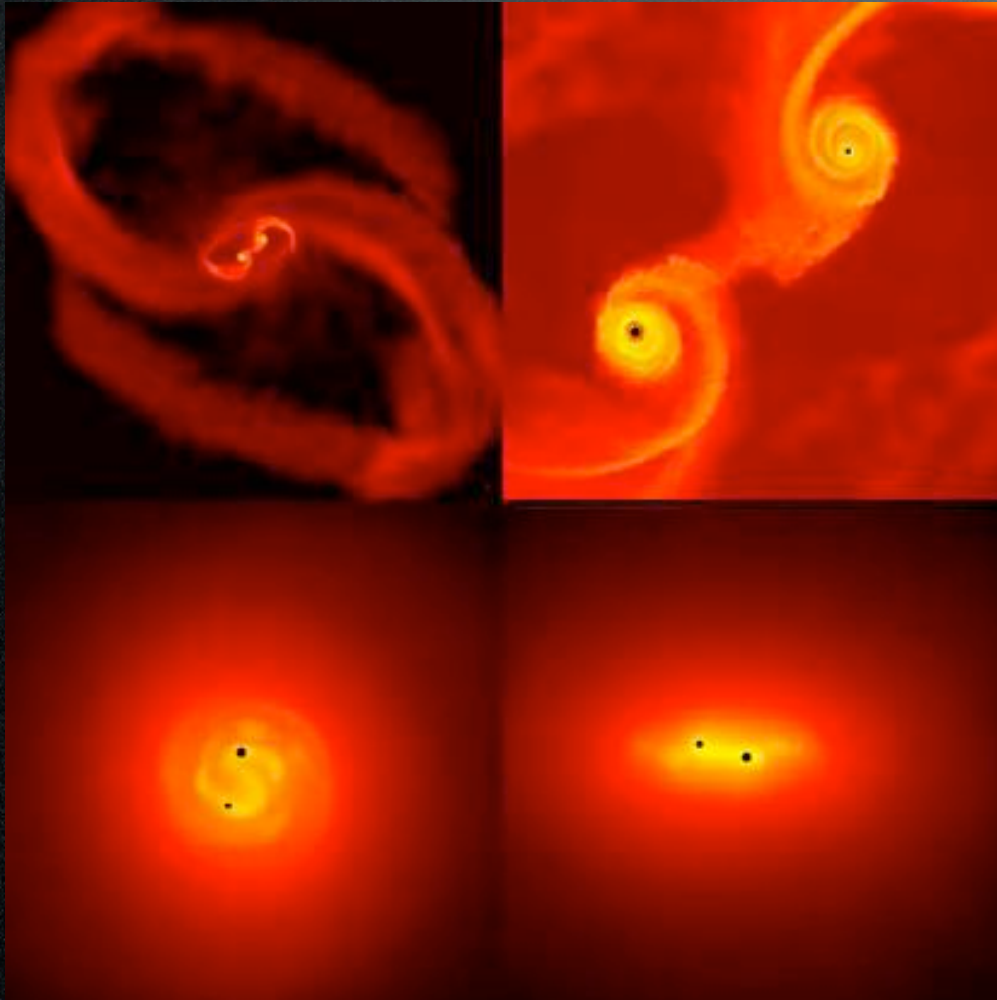
Gas accretion vs BH-BH mergers

Can we meet the constraints at low-ish redshift?

Matching the LF of quasars

How does the SMBHs mass grow along the cosmic history?

- **Mergers**
- **Accretion**



How does the SMBHs mass grow along the cosmic history?

→ Mergers

→ Accretion

**Soltan's
argument**

$$\rho_{\text{qso}(0)}^{\text{B}} = 2 \times 10^5 [0.1(1-\epsilon)/\epsilon] M_{\text{sun}} \text{Mpc}^{-3}$$

$$\rho_{\text{qso}(0)}^{\text{X}} = 2 \div 4 \times 10^5 [0.1(1-\epsilon)/\epsilon] M_{\text{sun}} \text{Mpc}^{-3}$$

$$\rho_{\text{SMBH}} = 2.5 \div 4.5 \times 10^5 M_{\text{sun}} \text{Mpc}^{-3}$$

Yu & Tremaine 2002, Elvis et al 2002, Merloni et al 2004

→ $\epsilon \geq 0.06$ @ $z < 5$: Kerr BHs?

→ the final mass of the SMBHs is dominated by accretion, with mergers playing a secondary role

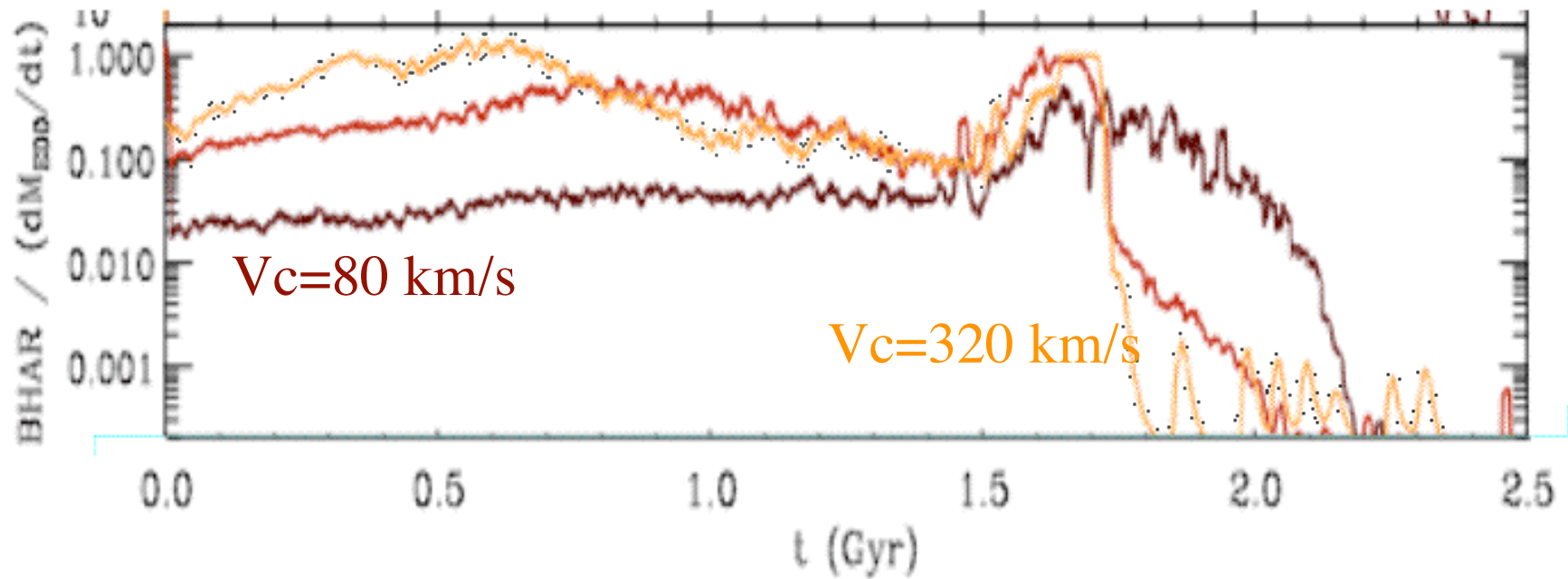
$$M_{\text{BH,acc}} \gg M_{\text{BH,merg}}$$

To recover the local M_{BH} - & the *quasar LF @ $z < 6$* :

✓ *only during major mergers* { *space density of quasars*
SMBHs-bulges connection

✓ *the accreted mass is a fixed fraction of the M_{BH} -relation* { *BH growth limited by feedback*
cfr. Di Matteo et al.

✓ *Eddington accretion rate* { *a sensible assumption?*

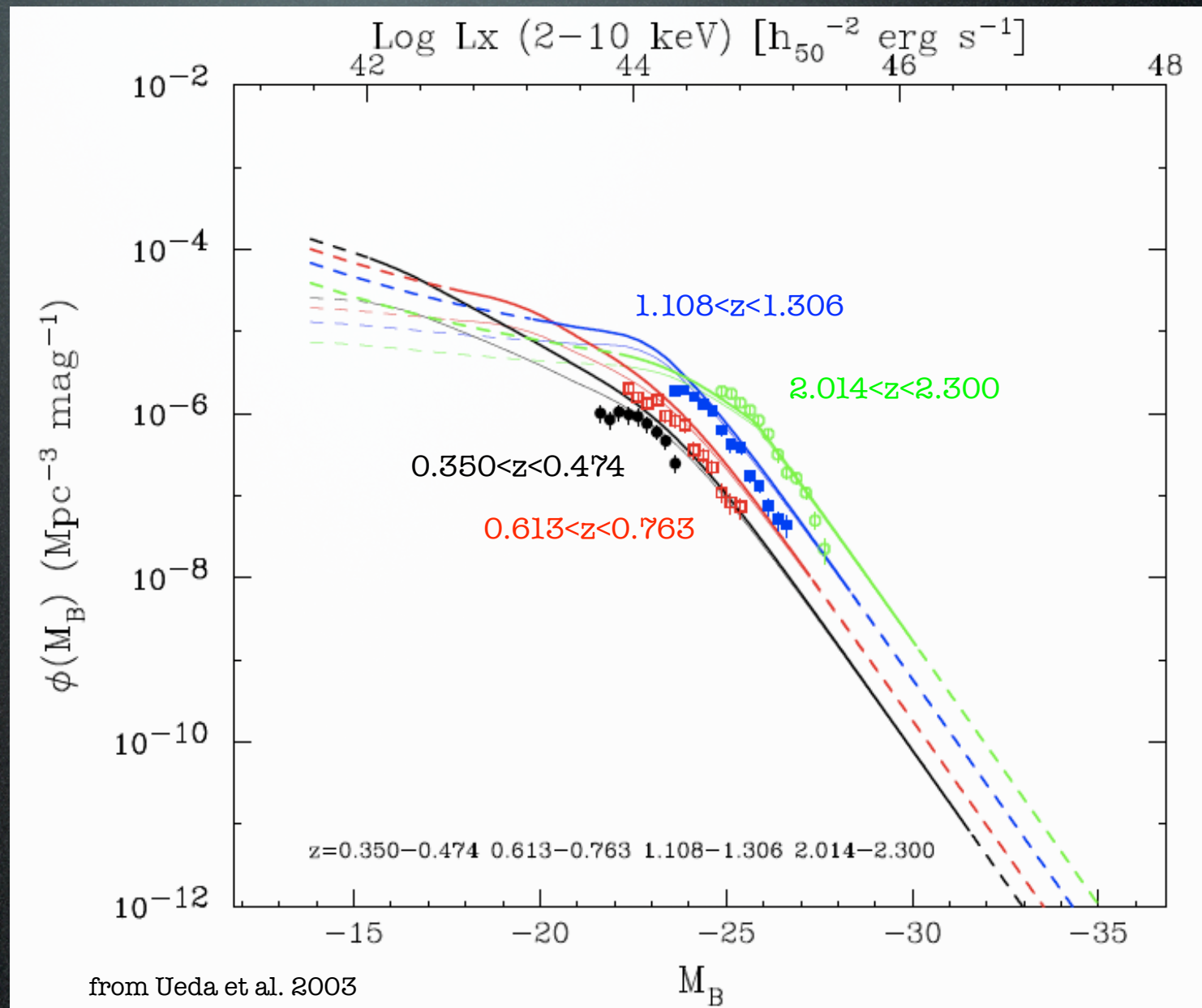


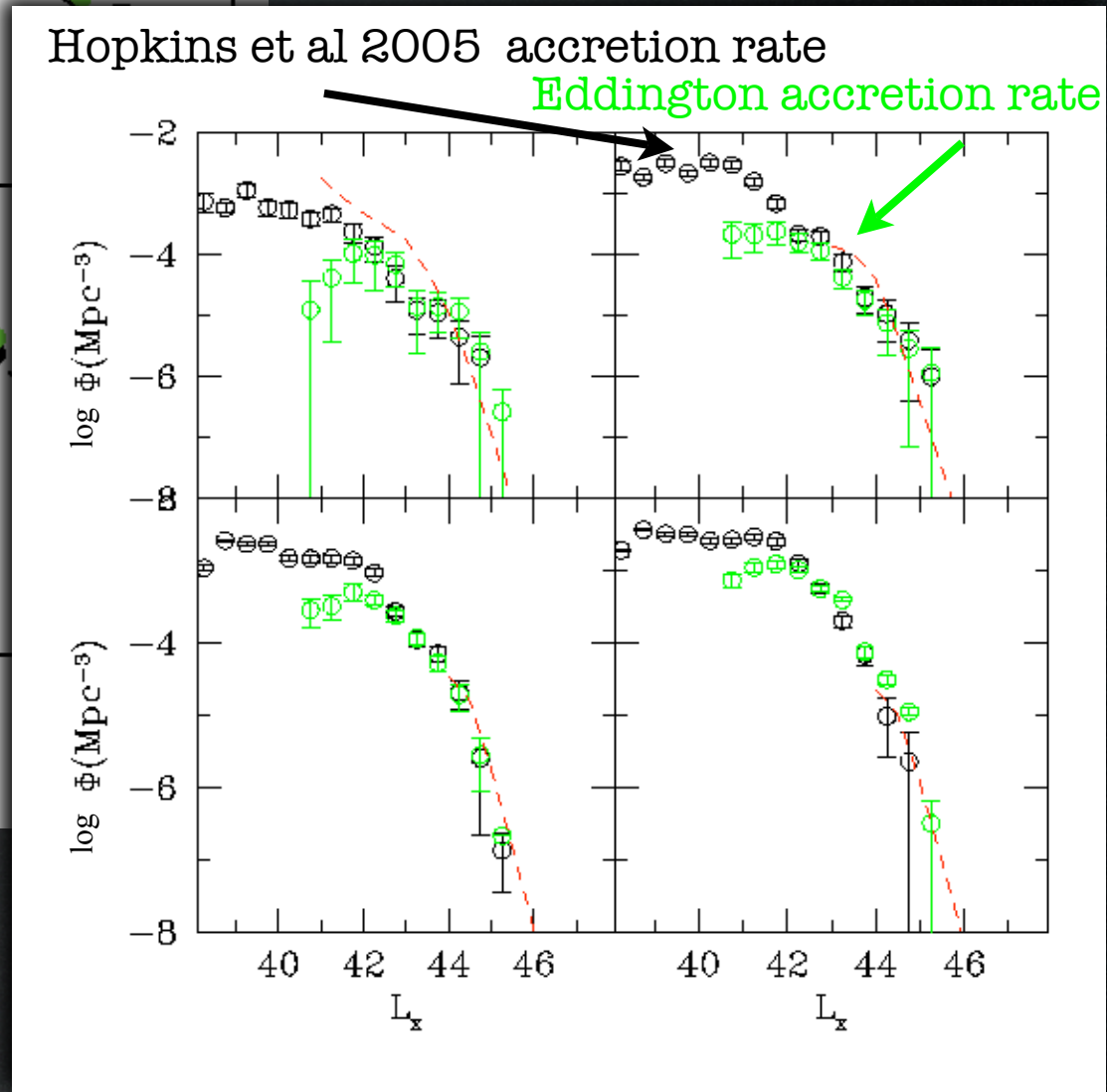
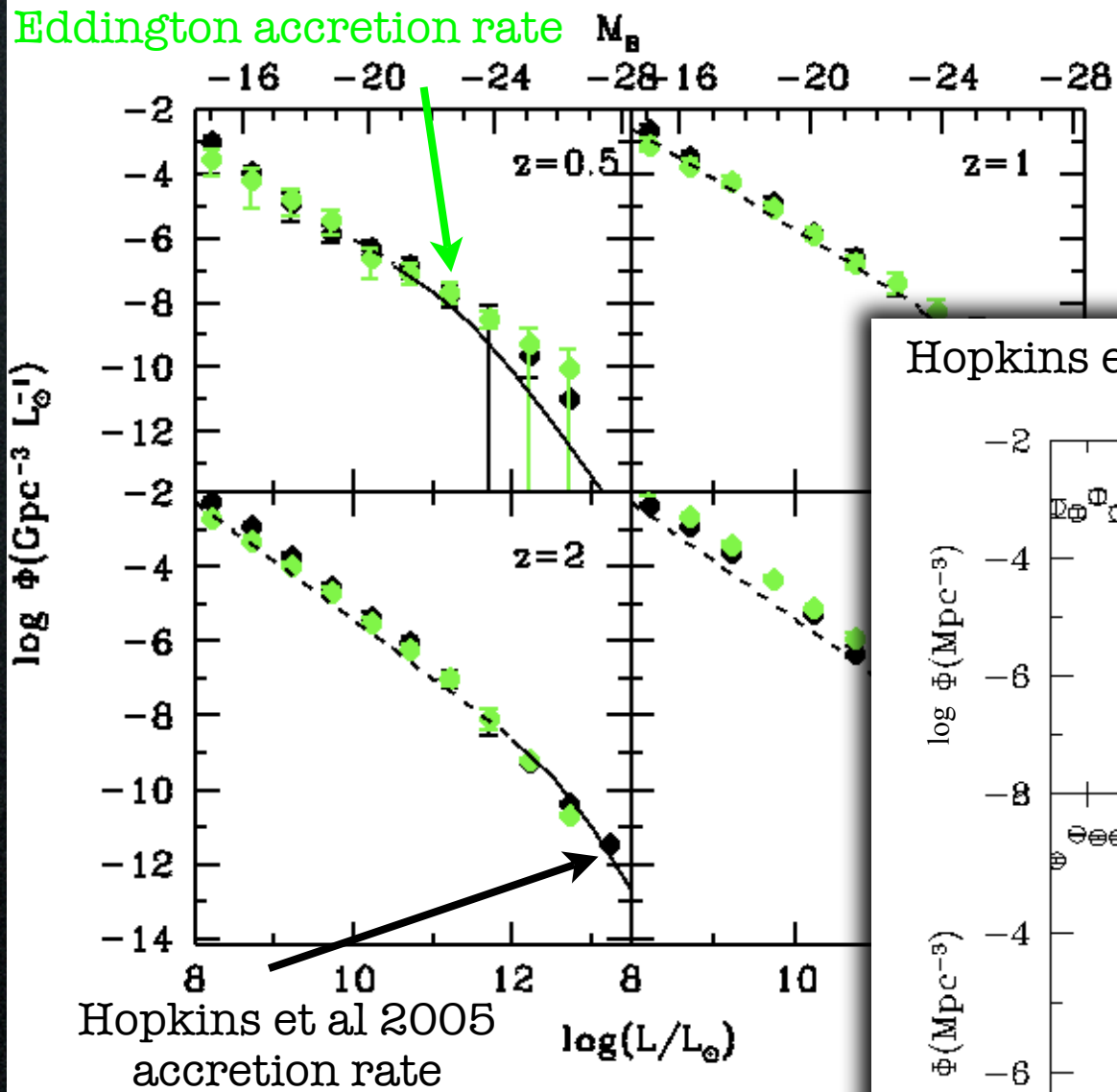
courtesy of T. di Matteo

Can the LF alone discriminate between models with a varying or constant Eddington ratio?

Testing the model at “low” redshift

Optical LFs
underestimate the
AGN population
BOTH at the **FAINT**
and **BRIGHT** end





Chapter 3

MBHs **weddings**...

MBH mergers: can we cut a long story short?

Most dynamical processes for MBH mergers imply long timescales. We observe very few systems... **nature must be much more efficient than our calculations...**

... and **honeymoons**

how MBHs take long trips, sometimes never to return

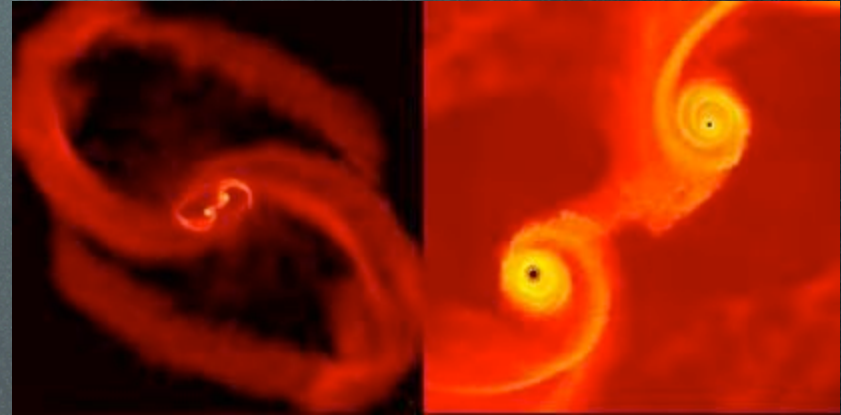
GAS DENSITY

Coplanar Equal-Mass Cooling+SF

10% Gas Fraction

time = 0.00 Gyr

Dynamical evolution of BH pairs



1. dynamical friction

✓ efficient only for *major mergers* against mass stripping

minor mergers: orbital decay > Hubble time

$$t_{\text{df}} \propto \frac{r_{\text{circ}}^2 V_c}{M_s} \propto \frac{1+P}{P} \frac{1}{H \sqrt{\Delta_{\text{vir}}}}$$

2. hardening of the binary (Quinlan 1996, Merritt 1999, Milosavljevic & Merritt 2001)

the binding energy of the BHs is larger than the thermal energy of the stars

$$a_h = \frac{Gm_2}{4\sigma^2} = 1 \text{ pc} \left(\frac{m_2}{2 \times 10^7 M_\odot} \right) \sigma_{150}^{-2}$$

X 3 bodies scattering between the binary and the surrounding stars: the SMBHs create a stellar density core ejecting the background stars

→ explain shallow cores in bright ellipticals

$$\frac{d\mathcal{M}_{ej}}{d \ln(1/a)} = J(m_1 + m_2)$$

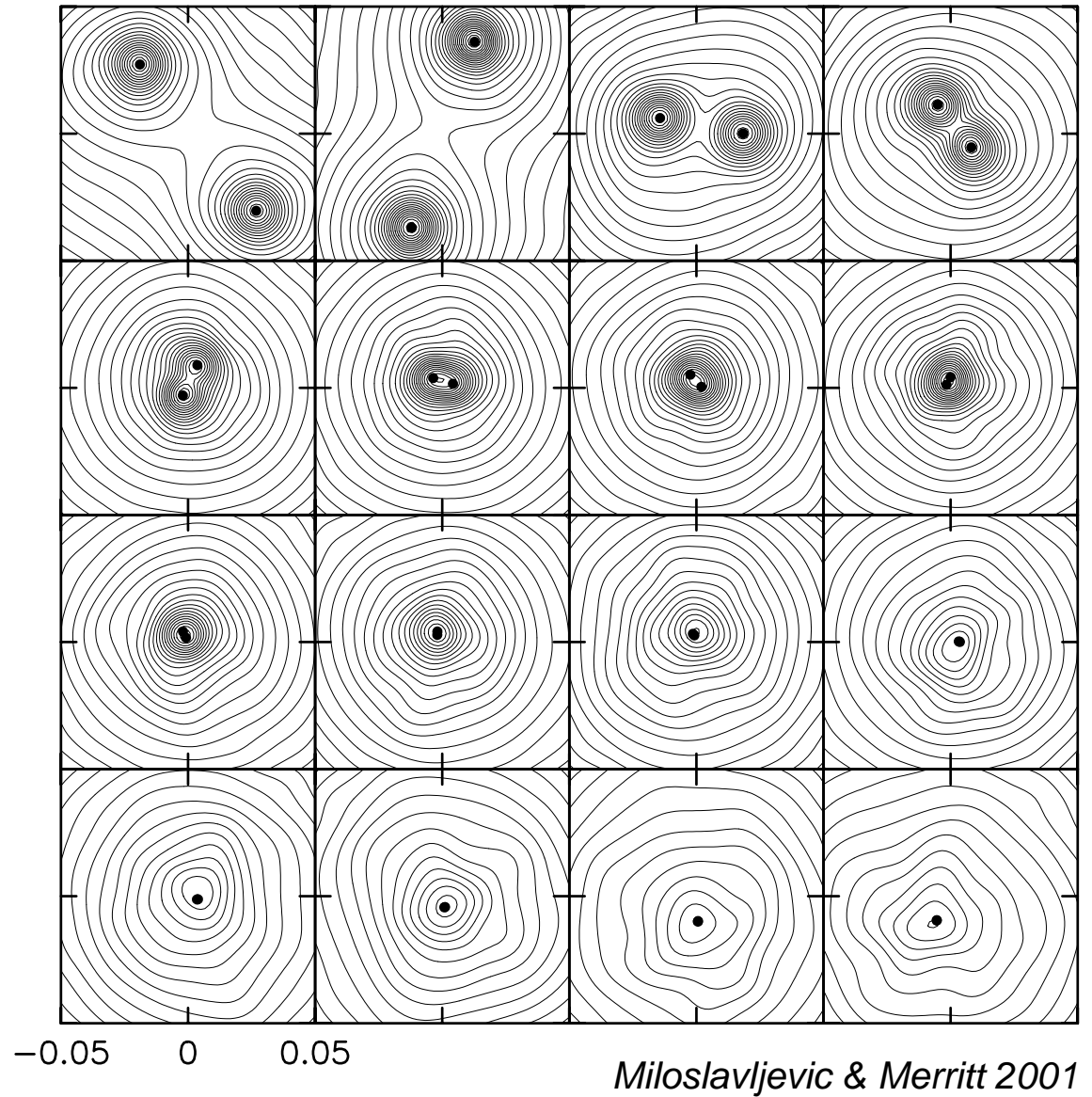
As the binary shrinks, ejecting a mass in stars comparable to its own mass, the star density decreases and the merging timescale increases

$$\rho_c(t) = \frac{8\sigma^6}{9\pi G^3 \mathcal{M}_{ej}^2(t)}$$

$$t_h = \frac{\sigma}{G\rho_* a H}$$

As the binary shrinks
ejecting stars the central
density drops

Numerical simulations do
not have the required
resolution (relaxation,
wandering...) yet?



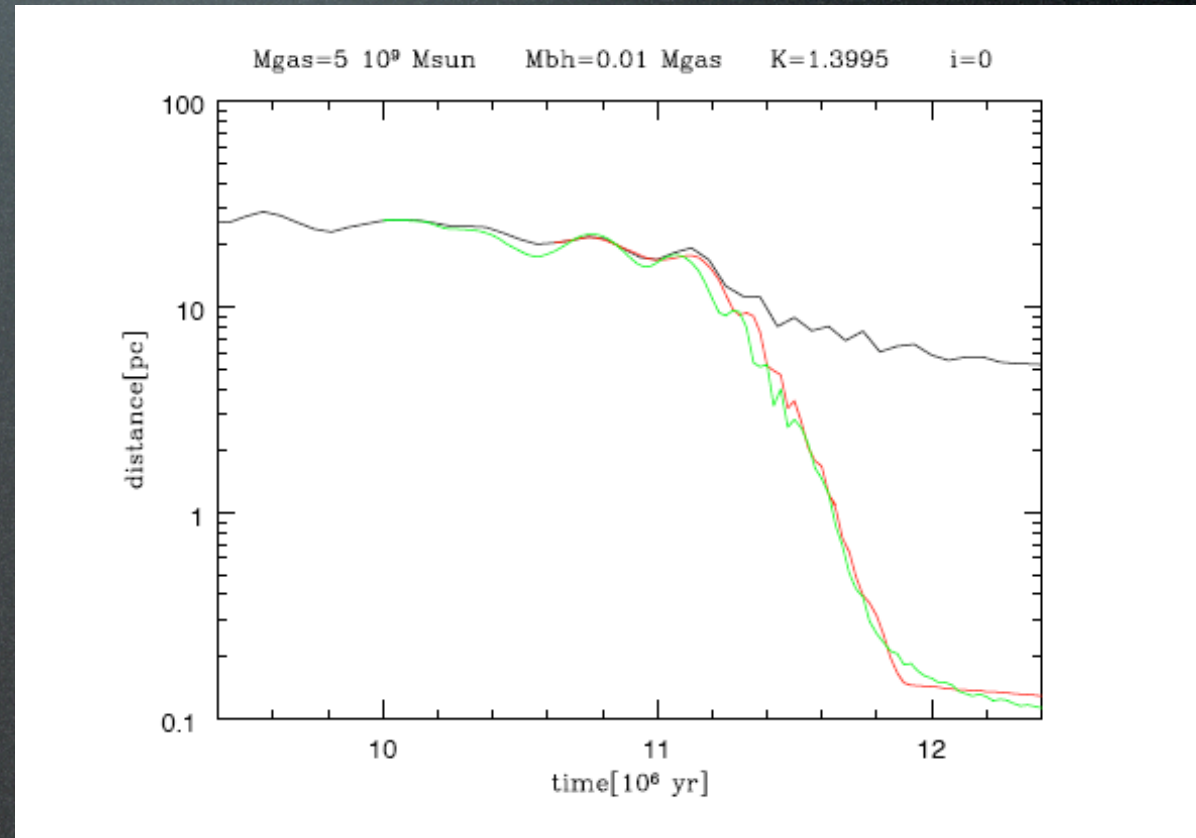
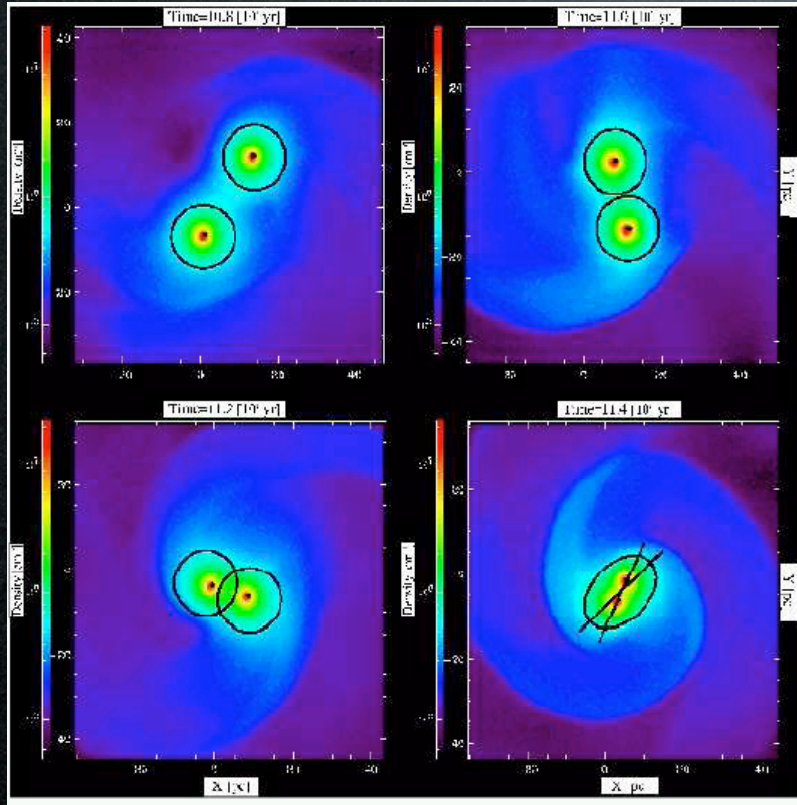
The BH sphere of influence is typically smaller than the core radii:
additional processes to enlarge the core (e.g. hierarchical evolution,
heating by reprocessed BHs orbital decay)

2. hardening of the binary (Quinlan 1996, Merritt 1999, Milosavljevic & Merritt 2001)

X *interaction with gas/accretion disc: much more efficient than stellar scatterings (Armitage & Natarayan, Escala et al)*

Gas is much more efficient than stellar scattering (Escala et al. 2004, Dotti et al 2006)

High-z galaxies: plenty of gas, few stars

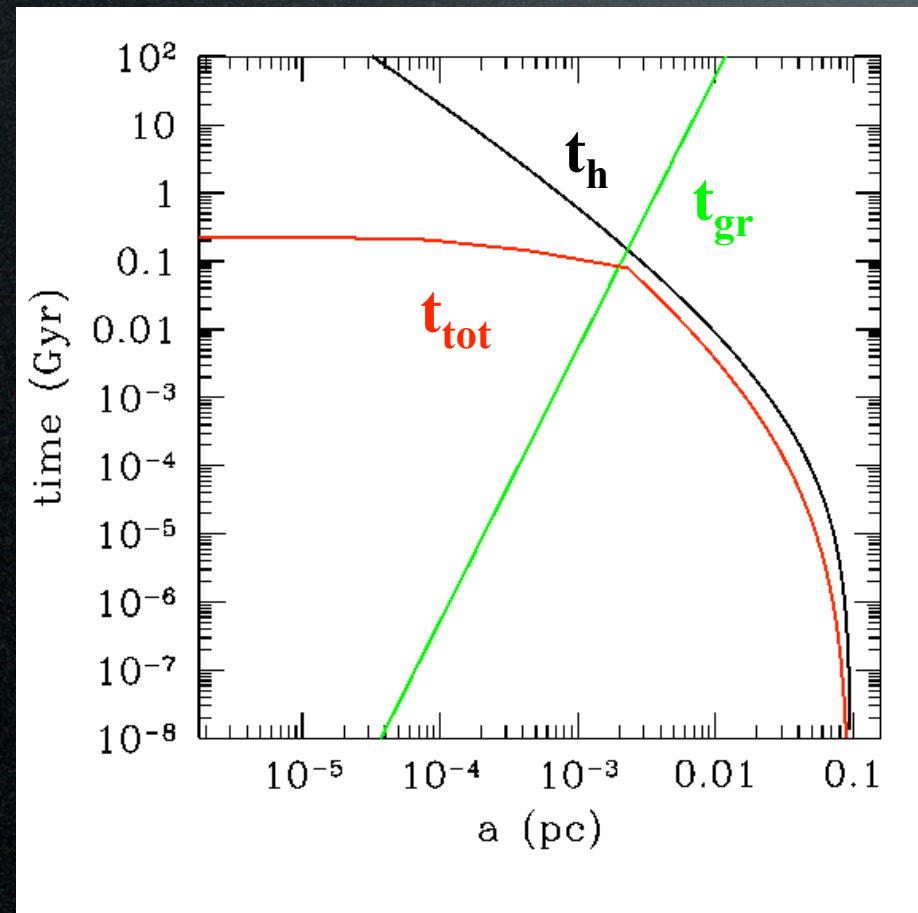
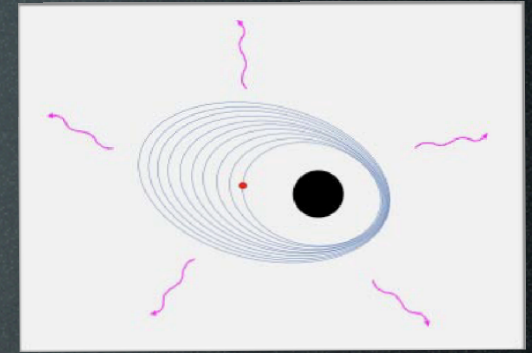


Also, if an **accretion disc** is present, the satellite BH is dragged in by **viscosity** on short-ish timescales (e.g. Armitage & Narayan, Ivanov et al.)

$$t_{\text{gr}}(0.1 \text{ pc}) \sim 200 \text{ Gyr}$$

3. emission of gravitational waves (Peters 1964)

Takes over at subparsec scales...



$$t_h = \frac{\sigma}{G\rho_0 a H} \quad \rho_c(t) = \frac{8\sigma^6}{9\pi G^3 M_{ej}^2(t)}$$

$$\frac{dM_{ej}}{d\ln(1/a)} = J(m_1 + m_2)$$

$$t_{gr} = \frac{5c^5 a^4}{256G^3 m_1 m_2 (m_1 + m_2)}$$

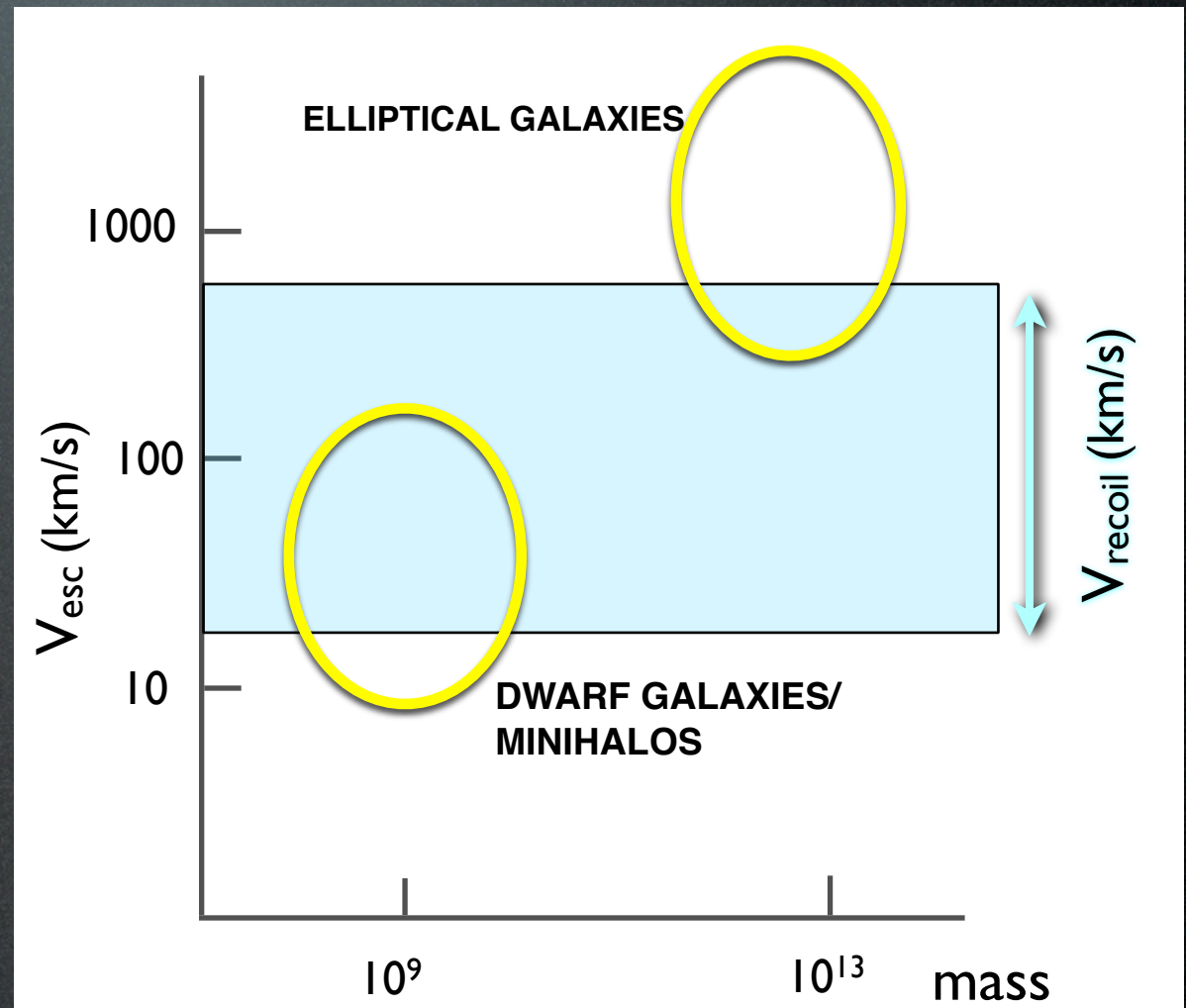
Gravitational Rocket

binary center of mass recoil during coalescence due to asymmetric emission of GW

(e.g. Fitchett 1983, Favata et al 2004, Blanchet et al 2005, Baket et al 2006)

**BH mergers:
positive or
negative
contribution?**

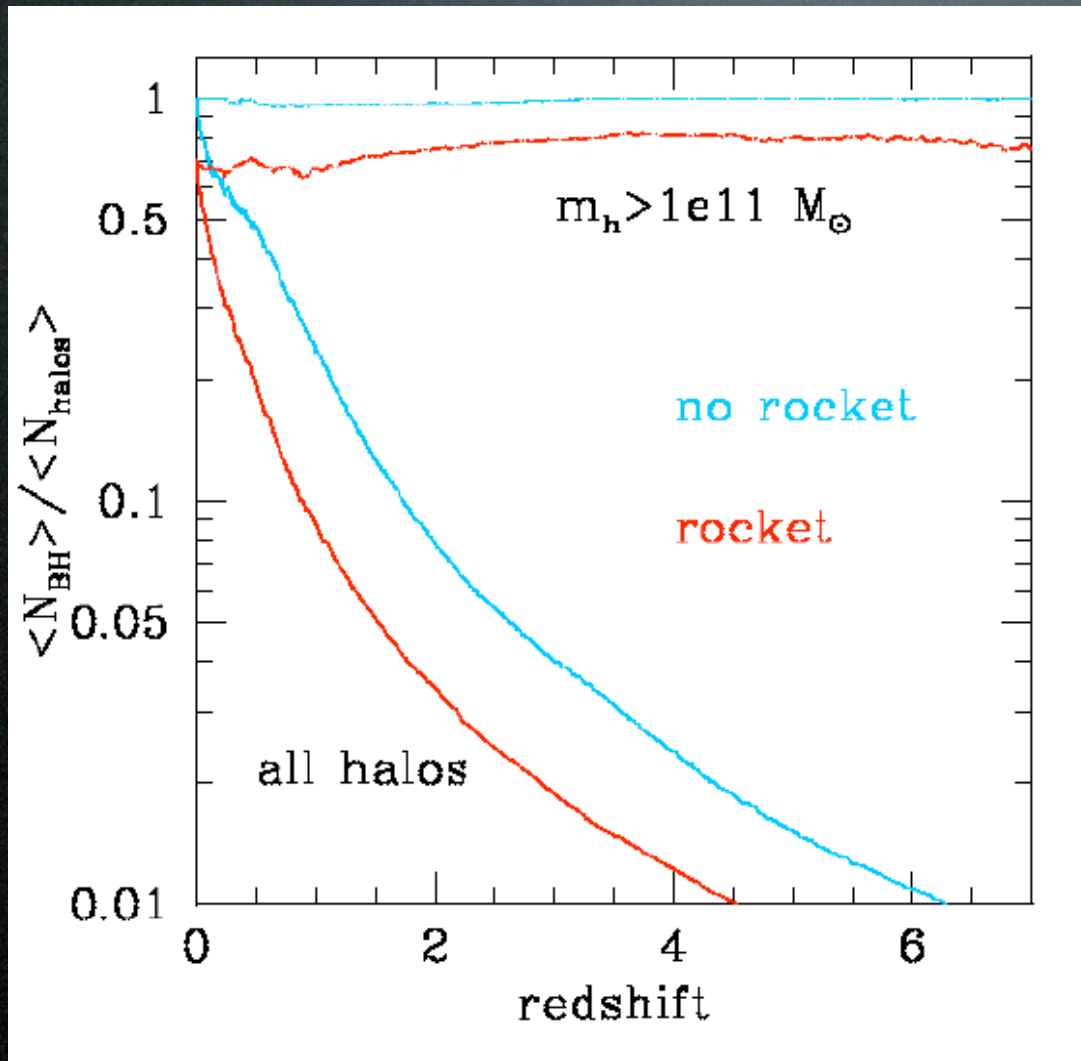
**Are merging BHs
ejected from
galaxies?**



*If mergers do happen **efficiently**, how much is the **gravitational rocket** effect a **threat** for the SMBH evolution?*

- 1. not ALL BHs experience a merger in their lifetime, only ~40-50%*
- 2. about 50% of the merging BHs are ejected out of their host halo, but 80% at $z > 10$!*
- 3. SMBHs gain their mass primarily by accretion at "low-ish" redshift ($z < 5$)*

MBH Occupation Fraction



for halos with mass larger than $3e12 M_{\text{sun}}$, the occupation fraction is unity, though

Chapter 4

We can model the mass growth of MBHs as traced by evolution of the quasar LF

The comparison between the total qso energy density and the SMBH mass density suggest a radiative efficiency $\epsilon \geq 0.1$

Are BHs rapidly spinning?

What is the typical BH spin predicted by the hierarchical evolution?

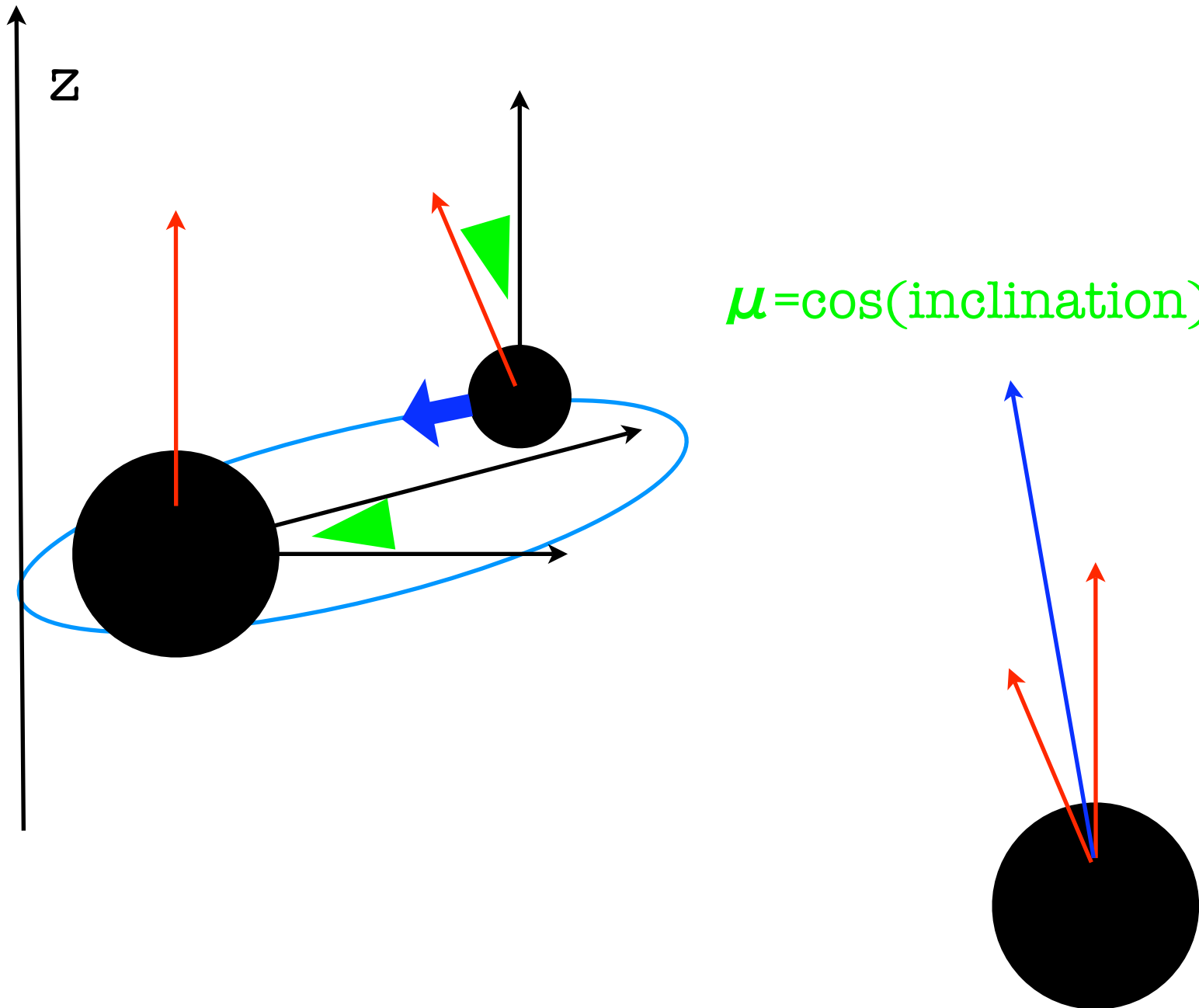
*What is the typical radiative efficiency value?
Are BHs rapidly spinning?*

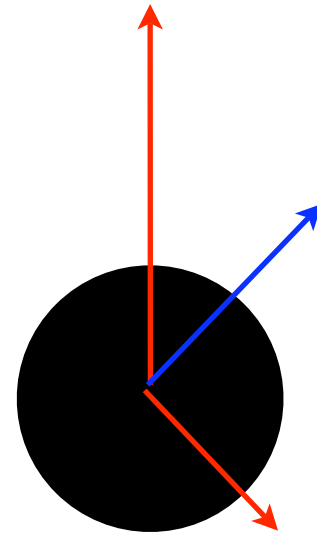
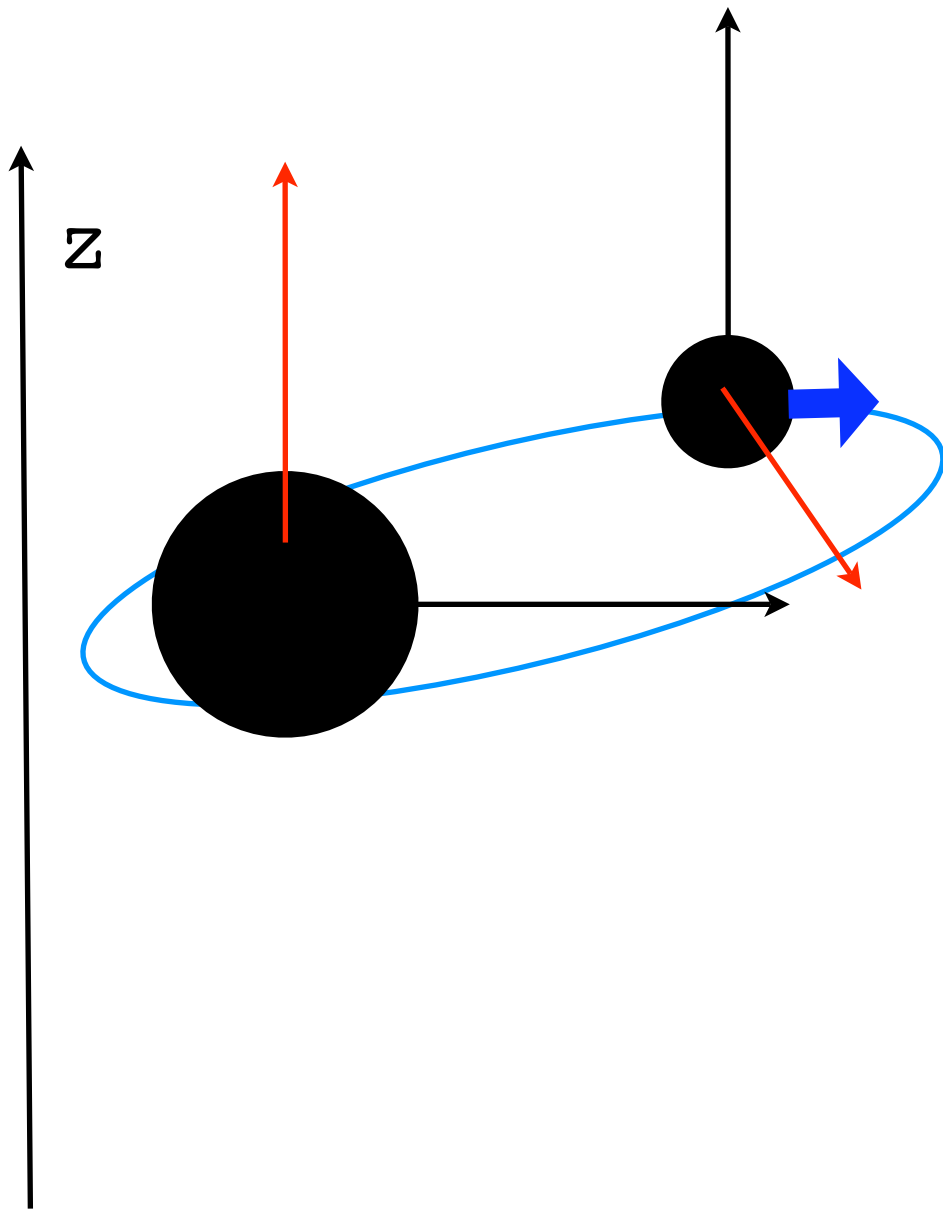
***BHs spin is modified by BH mergers and the
coupling with the accretion disc***

- ✓ mergers can spin BHs either up or down*
- ✓ alignment with a thin disc spins up*

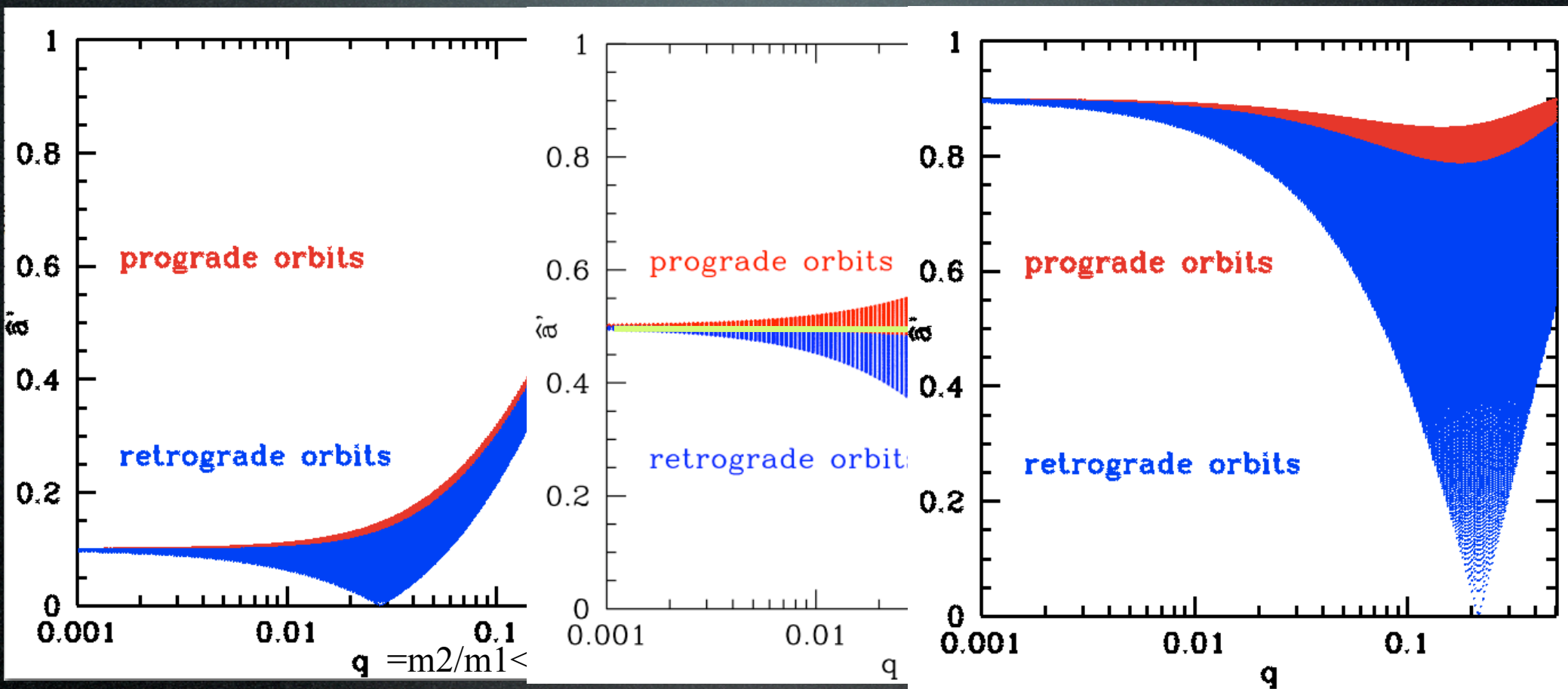
... but BH mergers are rare events!

Evolution of BH Spin: MERGERS





Evolution of BH Spin: MERGERS



Equatorial prograde/retrograde orbits

$$\hat{r}_{\text{LSO}} = 3 + Z_2 \mp \sqrt{(3 - Z_1)(3 + Z_1 + 2Z_2)}$$

$$\hat{L}_{\text{LSO}} = \pm \frac{2}{3\sqrt{3}} [1 + 2(3\hat{r}_{\text{LSO}} - 2)^{1/2}] \quad \hat{E}_{\text{LSO}} = \left(1 - \frac{2}{3\hat{r}_{\text{LSO}}}\right)^{1/2}$$

Orbits with random inclination: $\mu = \cos(\text{inclination})$

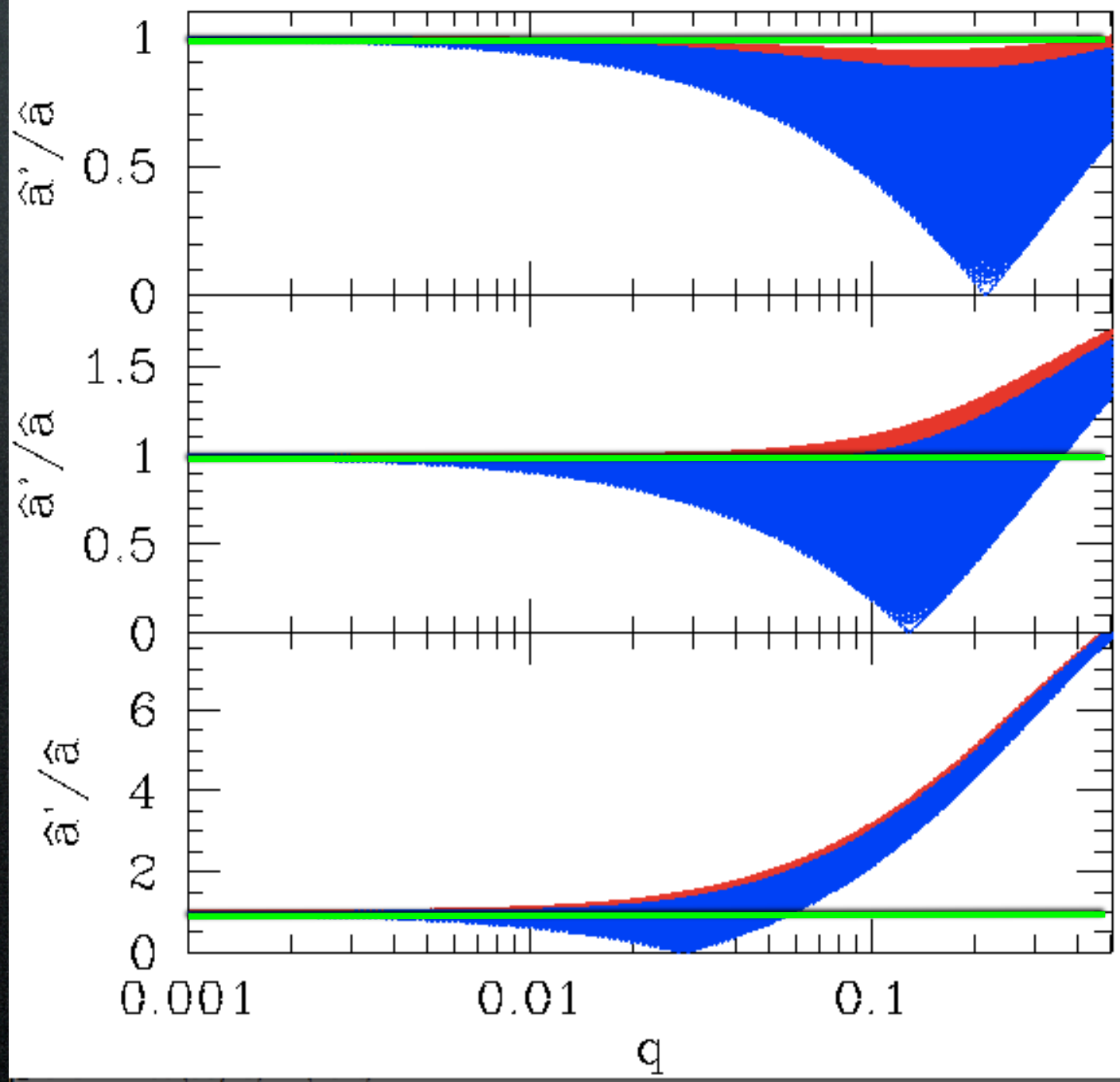
$$\hat{L}_{\text{LSO}} = \hat{L}_{\text{LSO}}(m\mu, \hat{a}_1)$$

$$\hat{E}_{\text{LSO}} = \hat{E}_{\text{LSO}}(m\mu, \hat{a}_1)$$

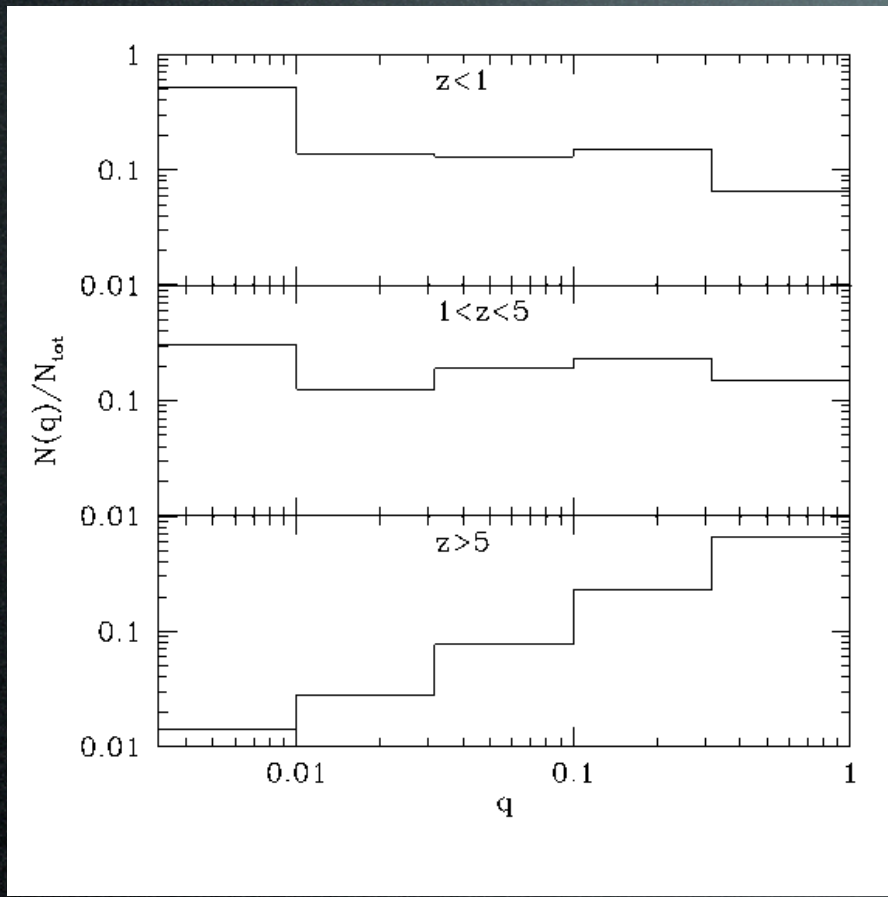
After the merger:

$$S'_z = m_1^2 [\hat{a}_1 + q\mu |\hat{L}(\hat{a}_1, \mu)|] \quad S'_\perp = qm_1^2 |\hat{L}(\hat{a}_1, \mu)| \sqrt{1 - \mu^2}$$

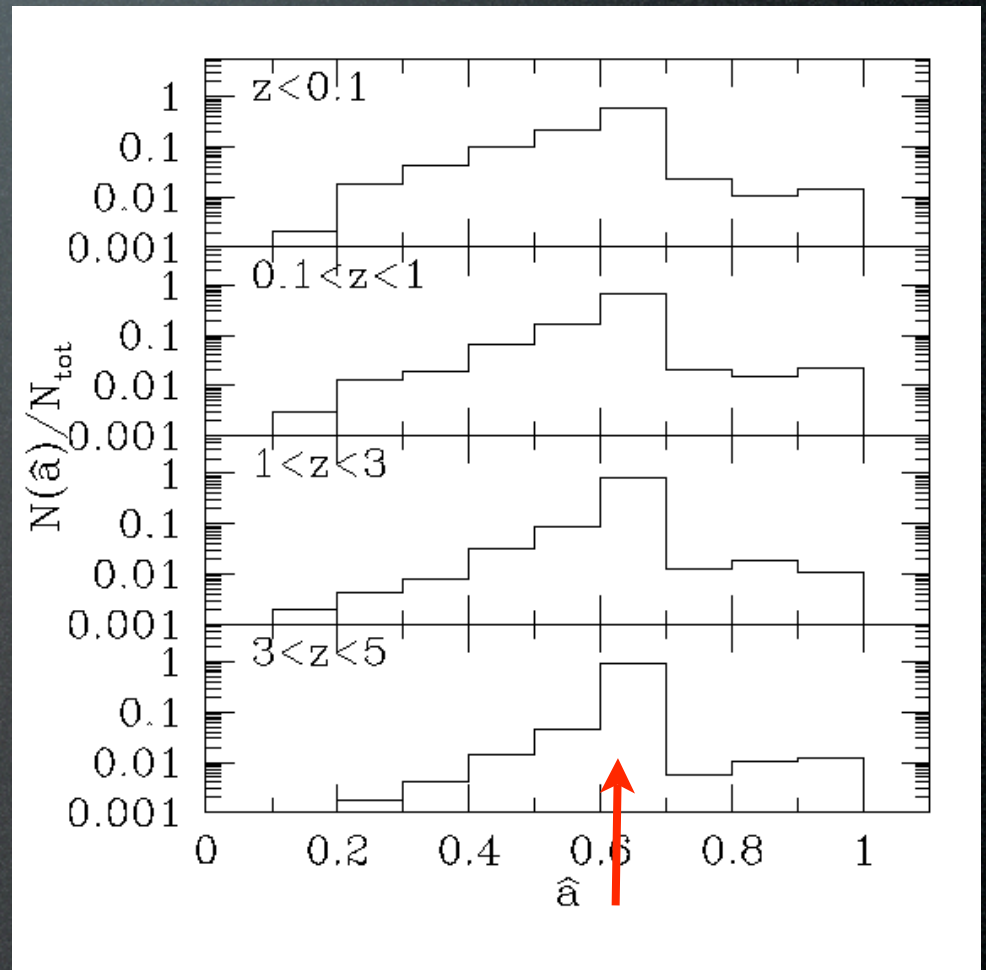
$$\hat{a}' \simeq \frac{\sqrt{S_z'^2 + S_\perp'^2}}{(1 + q)^2 m_1^2} \quad \hat{a}'^2 \simeq \frac{\hat{a}_1^2 + q|L|(q|L| + 2\hat{a}_1\mu)}{(1 + q)^4}$$



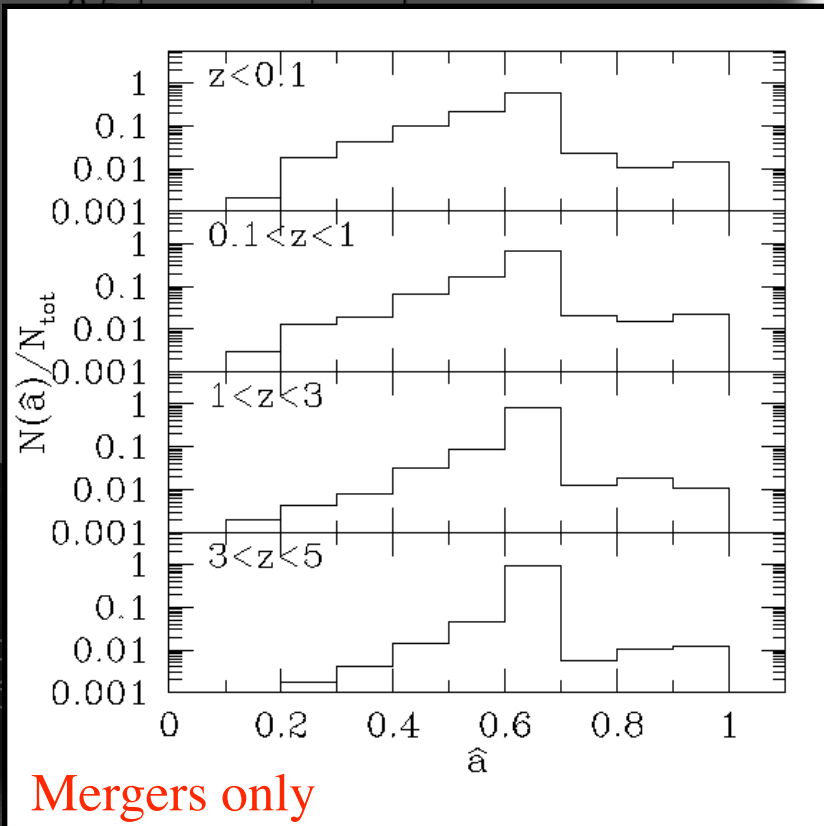
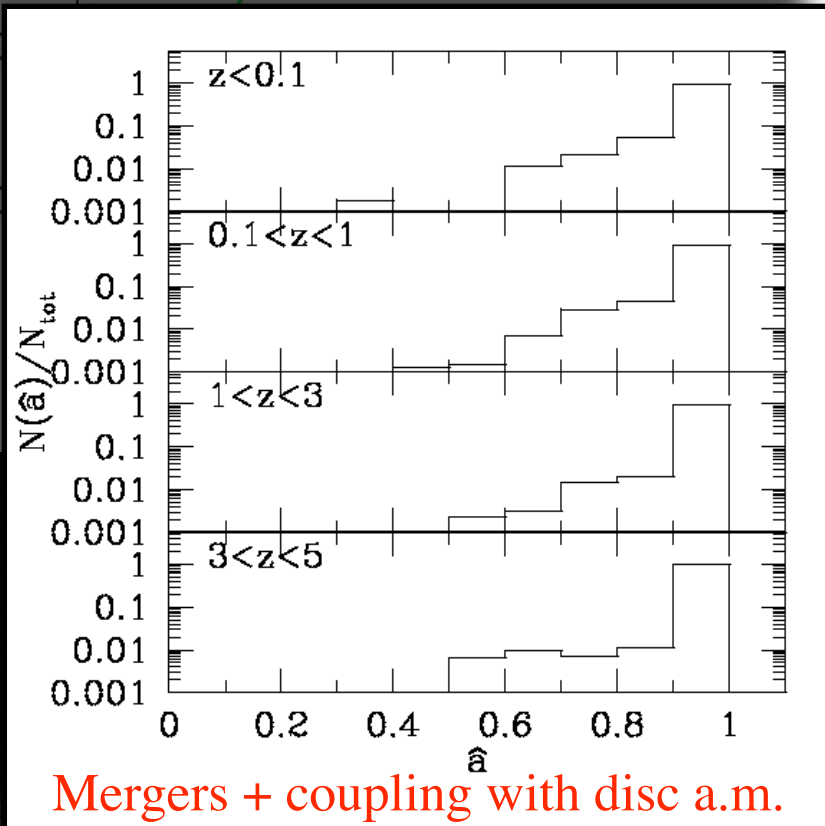
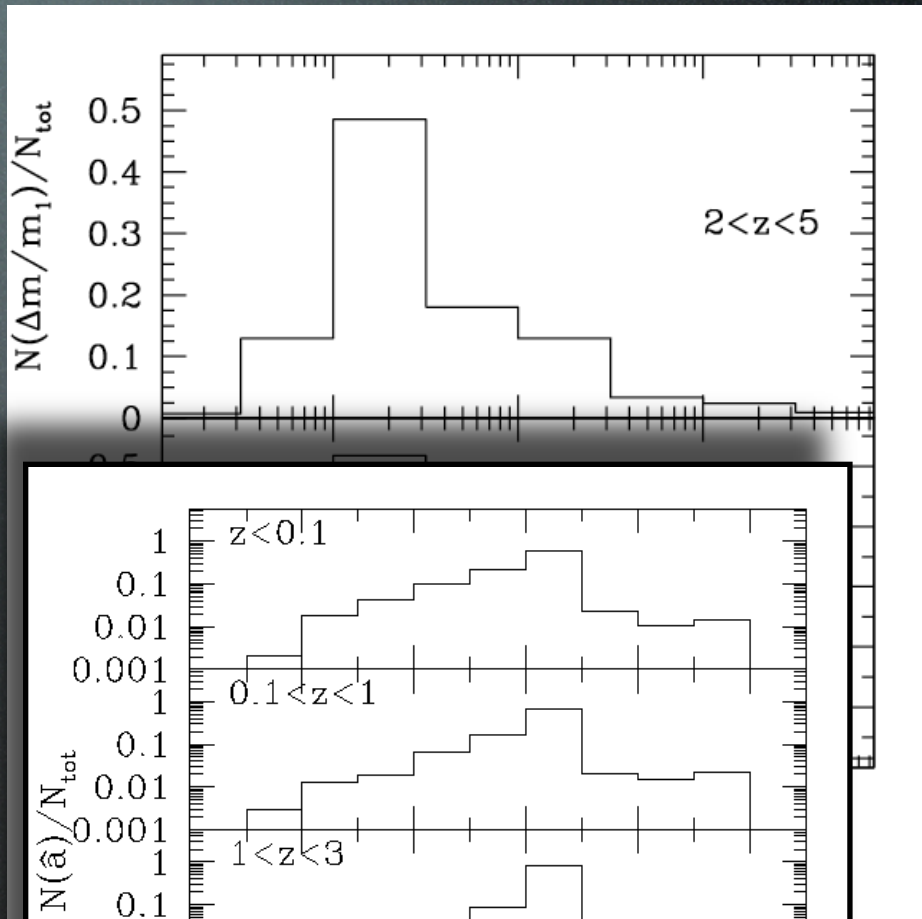
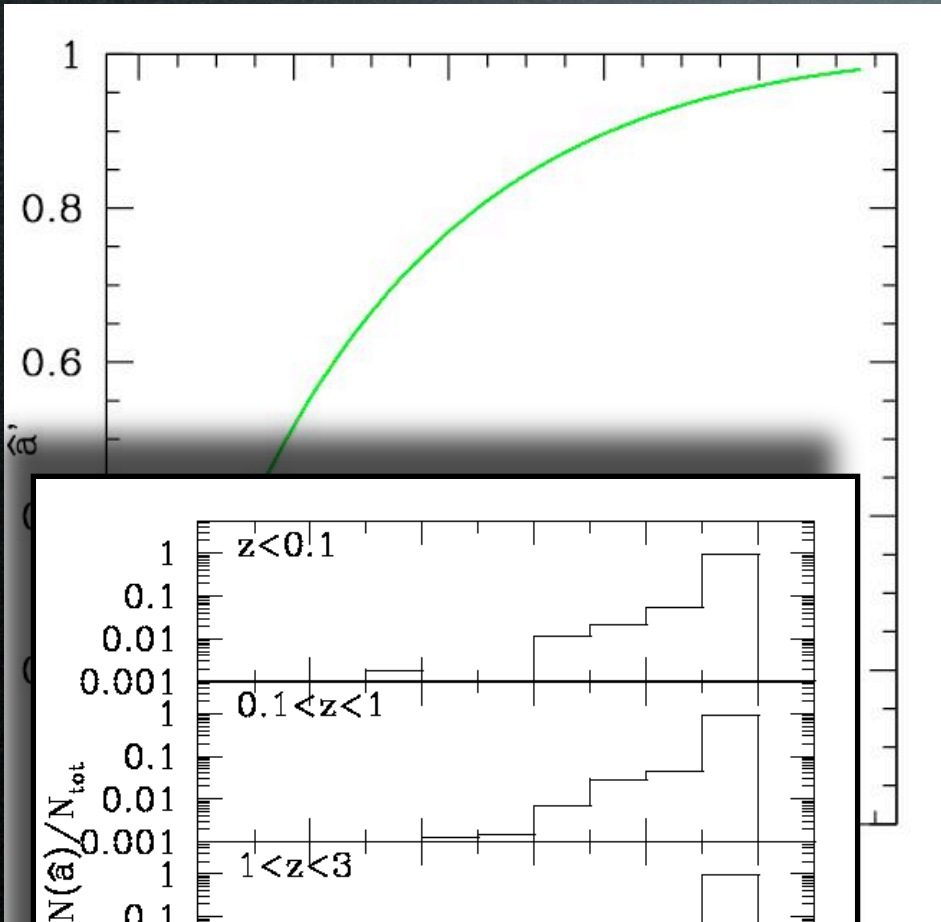
Distribution of black hole mass ratios from the whole cosm. evolution
 $q = m_2/m_1 \leq 1$



Distribution of black hole spins if only mergers influence them



Evolution of BH Spins



- ✓ via br
- ✓ the sof

et al 2005)

Mergers + coupling with disc a.m.

Mergers only

Chapter 6

How can we probe SMBH seeds? GW

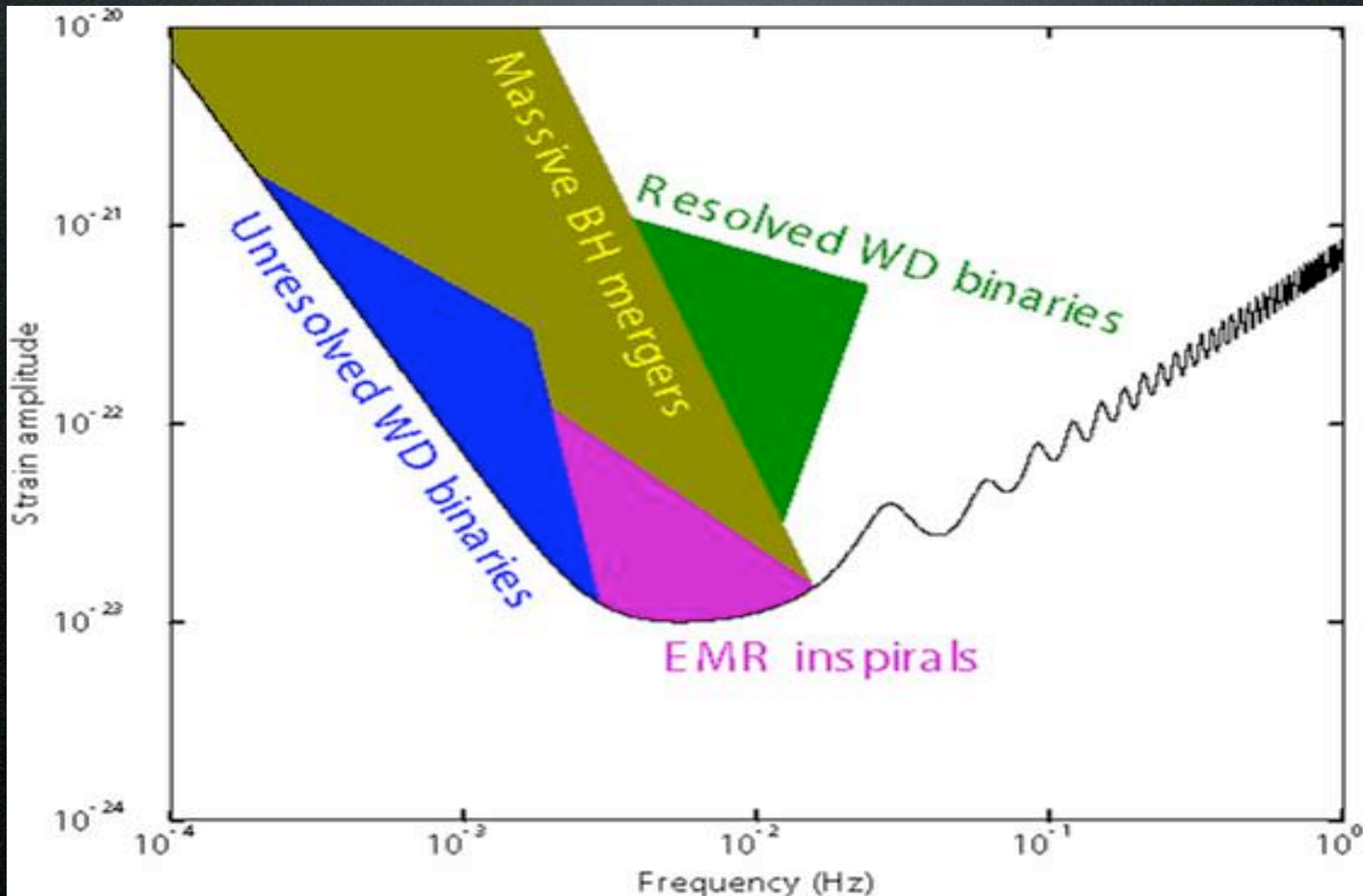
How can we probe SMBH spins? GW

How can we win a Nobel prize? GW ;)

... well if NASA behaves...

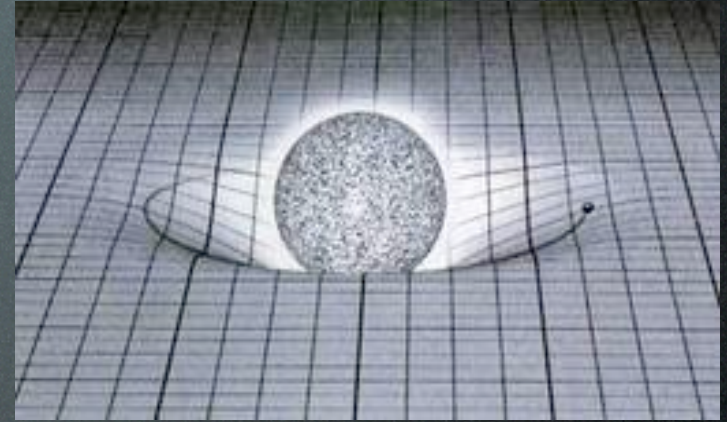
GWs from MBH binaries

Fluctuations in spacetime curvature, generated by rapidly accelerating masses: an exciting new window on the Universe to complement electromagnetic observations.

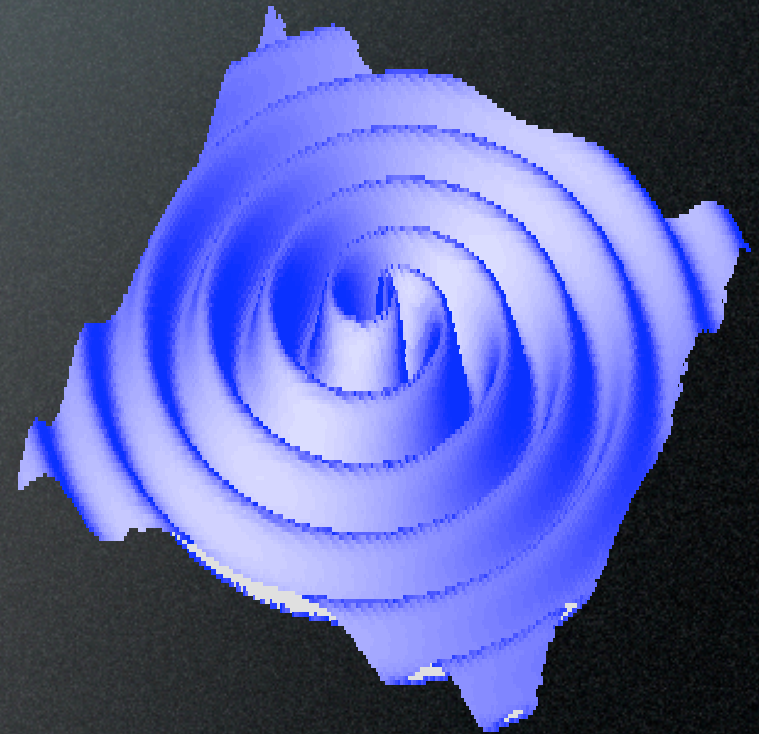


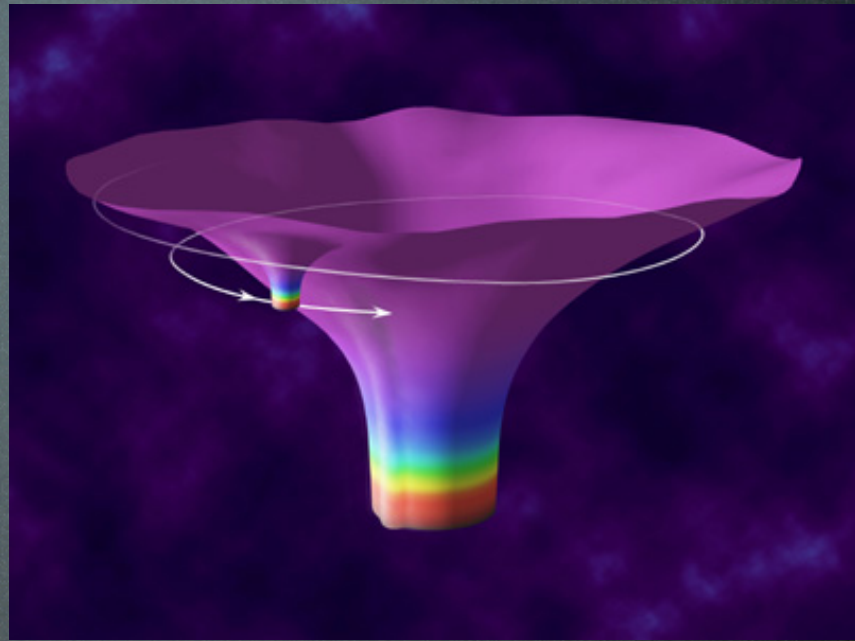
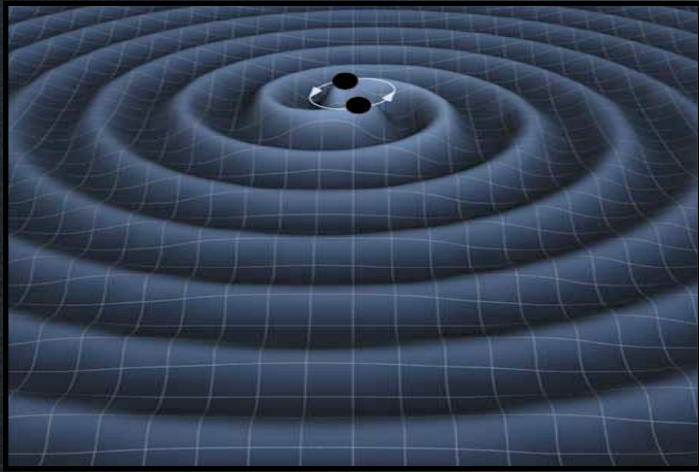
Gravity is spacetime curvature.

Any mass/energy bends spacetime near it.



Rapidly moving masses generate fluctuations in spacetime curvature: gravitational waves!





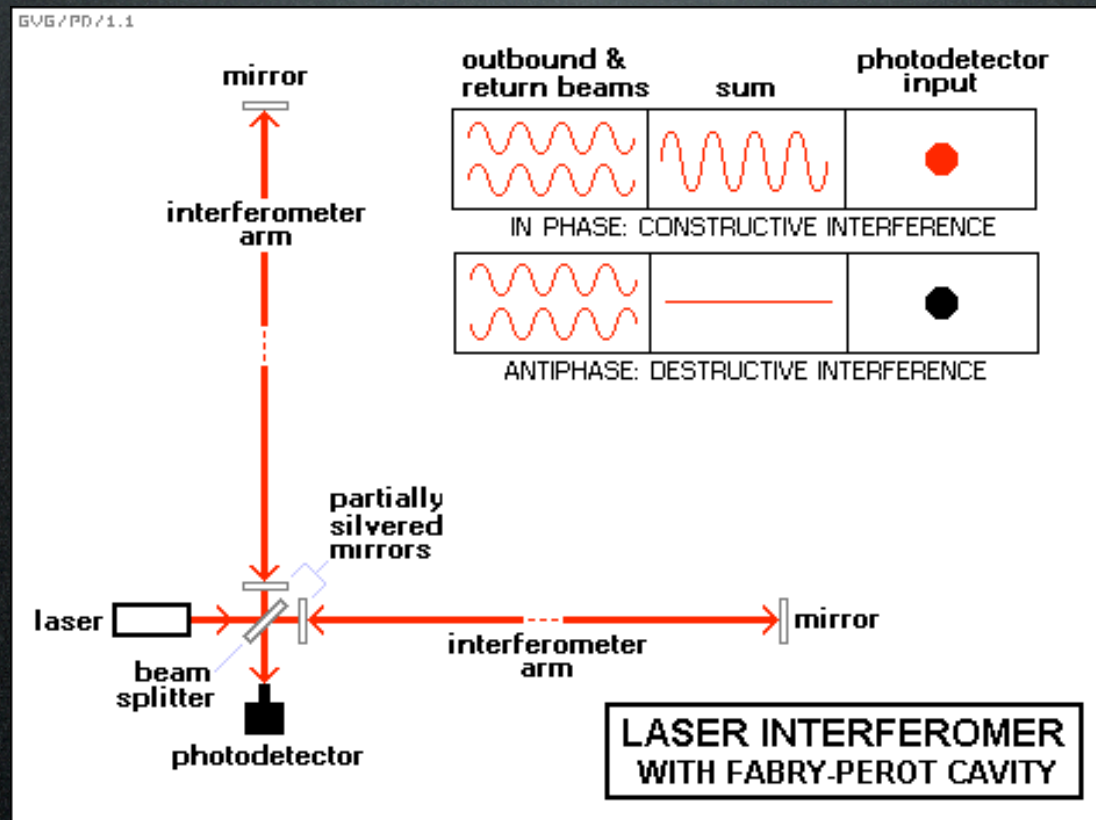
$$f = \frac{2f_{kepler}}{1+z} = \frac{1}{\pi(1+z)} \left(\frac{G(M_1 + M_2)}{r^3} \right)^{1/2}$$

$$r > R_{Sch} = \frac{2G(M_1 + M_2)}{c^2}$$

$$f < \frac{4 \times 10^{-3}}{1+z} \text{Hz} \left(\frac{10^6 M_{\odot}}{M_1 + M_2} \right)$$

Several **ground based interferometers** are now operating or are being built

Exploit quadrupole nature of GWs – space is distorted in opposite sense in two perpendicular directions – use a Michelson interferometer



Several **ground based interferometers** are now operating or are being built

$$L \propto \lambda \propto 1/f$$

$$h = \Delta L / L \sim 10^{-20} \quad \Delta L = 10^{-15} \text{ cm !!!}$$



LIGO – US project. Two 4km and one 2km detector

GEO – British/German project. One 600m detector

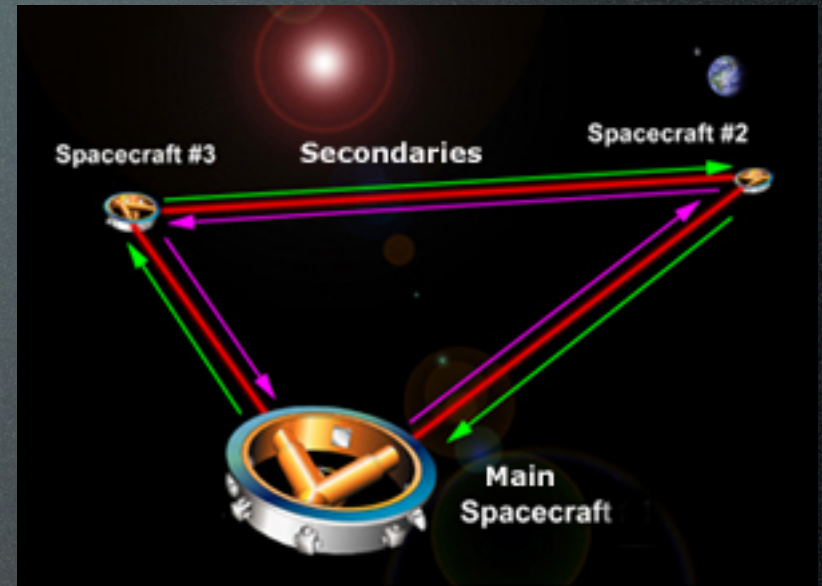
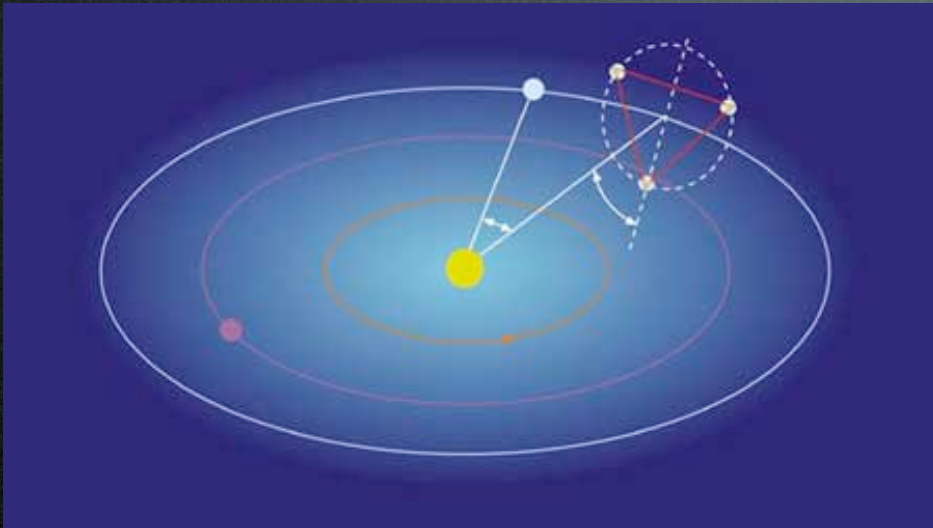
VIRGO – Italian/French project. One 3km detector

TAMA – Japanese project. One 300m detector

AIGO – Australian project. One 80m detector

space based interferometer: **LISA**

Joint NASA/ESA mission - launch date is 20??



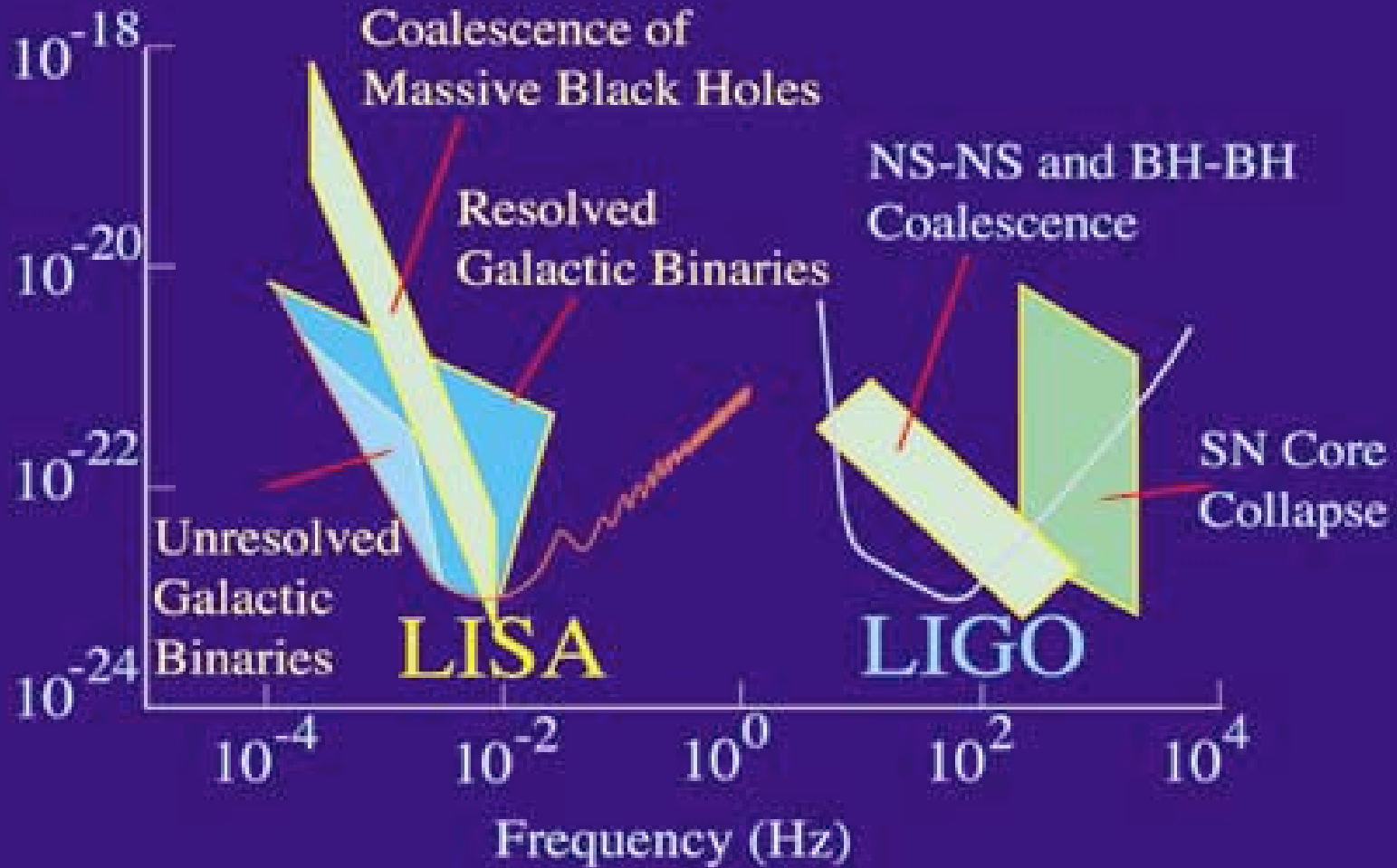
Longer baseline (5 million km) gives sensitivity to lower frequency gravitational waves

$$L \propto \lambda \propto 1/f$$

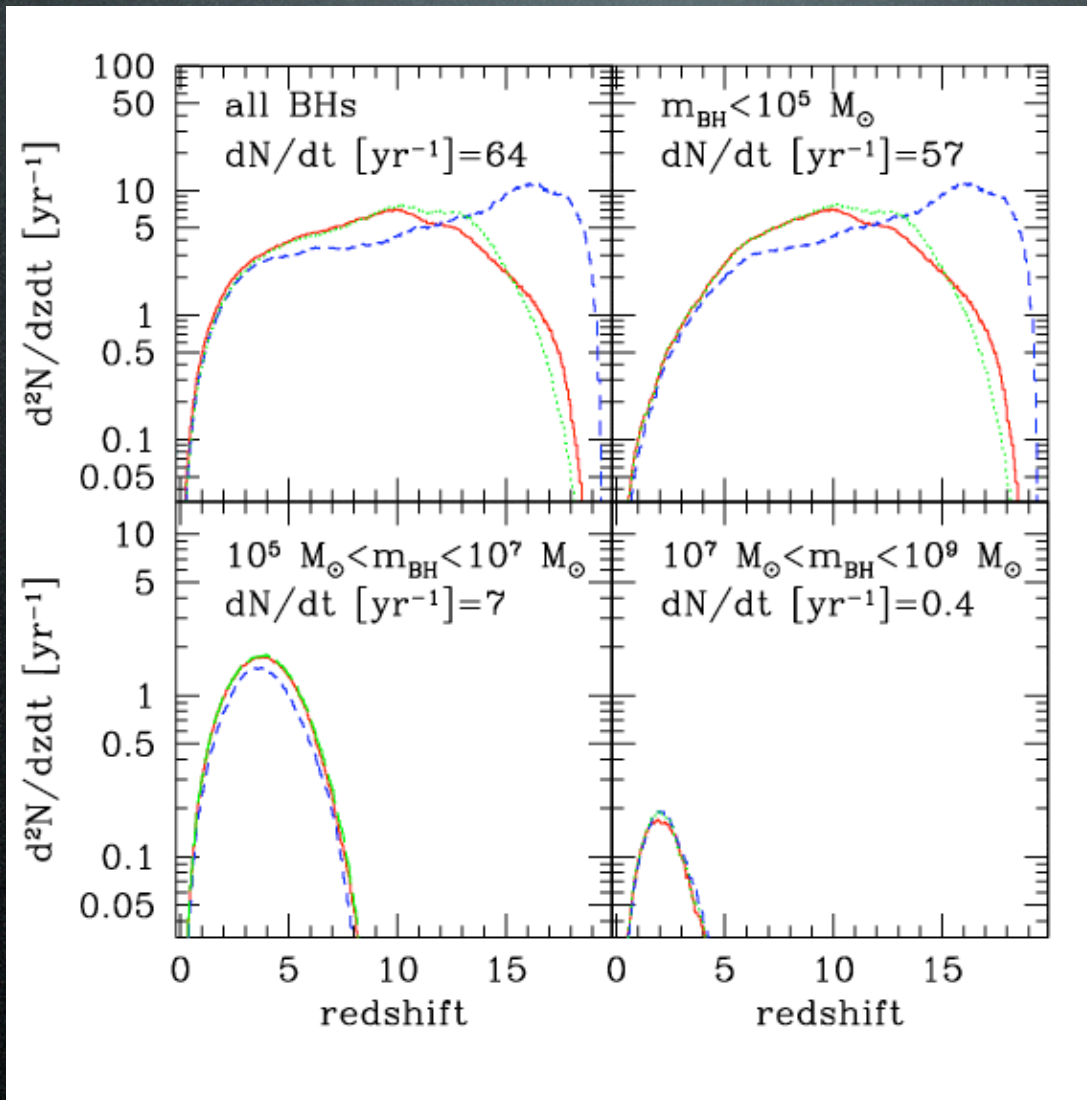
$$h = \Delta L / L \sim 10^{-20} \quad \Delta L = 10^{-8} \text{ cm for 5 million km!!!}$$

LISA will be a true GW telescope

Gravitational Wave Amplitude



SMBH binary merger rate



Emission spectrum of a MBH binary

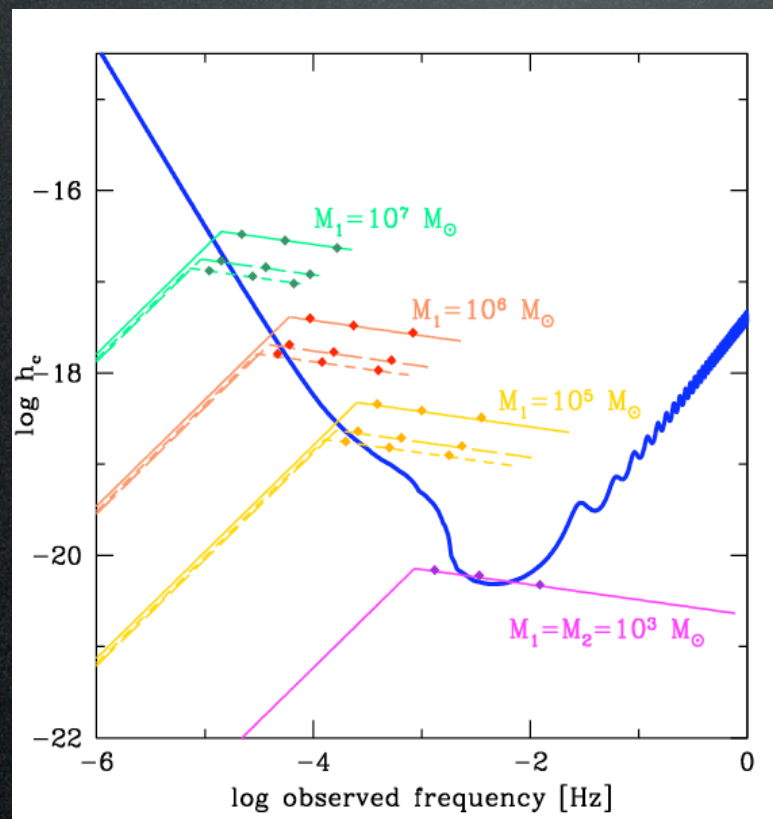
$$h_c^2(f) \propto \left(\frac{dE_{gw}}{df} f \right) f^{-2}$$

$$\frac{dE_{gw}}{df} f = \frac{dE_{gw}}{dt} \frac{dt}{dr} \frac{dr}{df} f$$

$$\frac{dr}{df} f = -\frac{2}{3} \left[\frac{G(M_1 + M_2)}{\pi^2} \right]^{1/3} f^{-2/3}$$

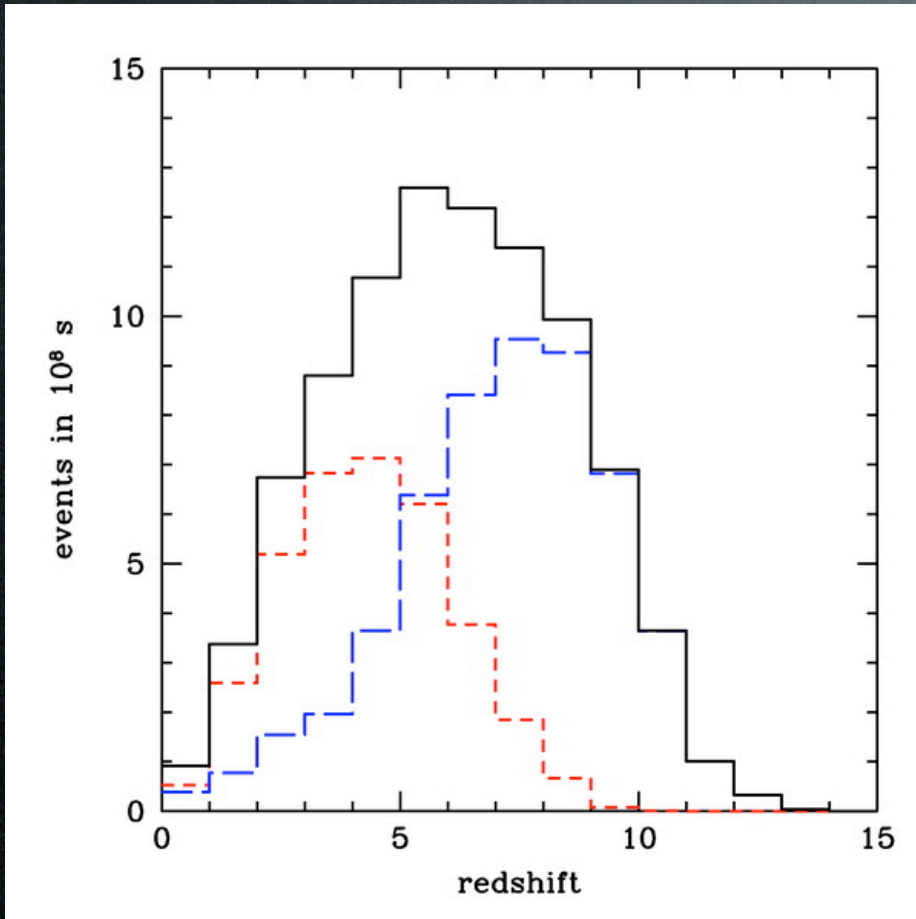
$$L_{gw} = \frac{32 G^{7/3} \pi^{10/3}}{5 c^5} (\mathcal{M} f)^{10/3}$$

$$h_c \propto f^{-1/6}$$



MBHs $M < 10^5 M_{\text{sun}}$
 can be detected up
 to $z=15-20$.
 Can really
 distinguish seed
 models

Resolvable events in 3 years



Typical masses $\sim 10^3$ - $10^6 M_{\text{sun}}$

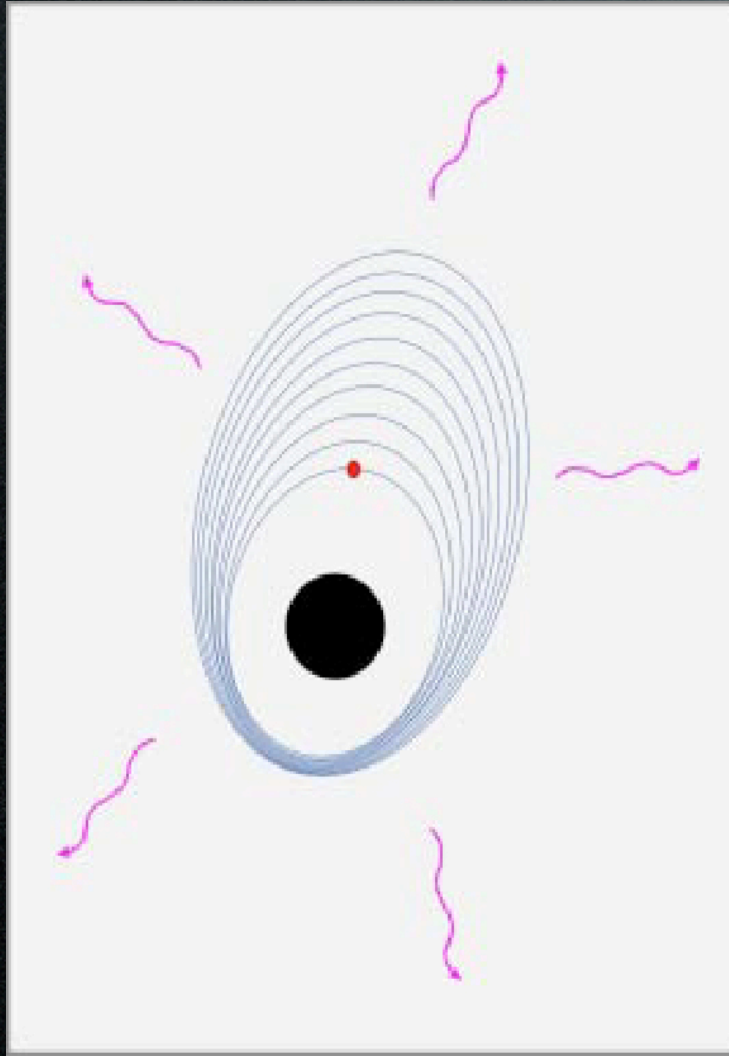
Typical mass ratio ~ 0.1

Sesana et al. 2004, 2005

Inspirals: longlasting wide binaries, small frequency change

Bursts: binaries coalescing during the observation period

Extreme mass ratio inspirals



Inspiral of a compact object (WD, NS, BH) into a supermassive black hole in the centre of a galaxy.

LISA can see $10M_{\text{sun}} + 10^6M_{\text{sun}}$ inspiral out to $z \sim 2 \rightarrow$ can probe SMBH spin evolution if event rate is high enough!!

For a typical event with $\text{SNR} \sim 30$, determine parameters with errors (Barack & Cutler, Creighton et al.)

$$\begin{aligned}M &\sim 2 \times 10^{-4} \\(S/M^2) &\sim 10^{-4} \\(\ln m) &\sim 10^{-4} \\(\ln D) &\sim 0.05\end{aligned}$$

Summary

SMBHs can be built up from seeds dating back to the end of the cosmological dark ages

- *seed BHs from the first stars, in the highest peaks of density fluctuations at $z \sim 20-30$*
- *mass accretion onto BHs triggered by major mergers efficiently spins-up black holes*
- *BHs interactions can displace BHs BUT they don't hinder the assembly of low- z SMBHs*
- *$z=6$ SDSS quasars require a large and quick mass build up: the early evolution can proceed differently - and more rapidly - in the most massive metal free halos: supercritical accretion*
- *If accretion is limited to the Eddington value dynamical effects must be mild*